

How much of the inflaton potential do we see?

Wessel Valkenburg
8 August, 2007

astro-ph/0703625, Phys.Rev.D75:123519, 2007

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Outline

Retrieving information on inflation

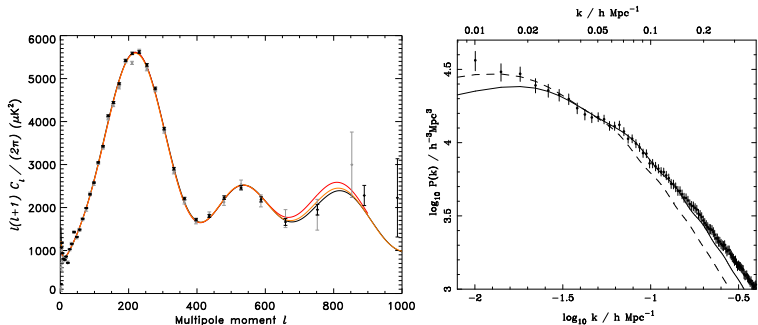
New results

Conclusion
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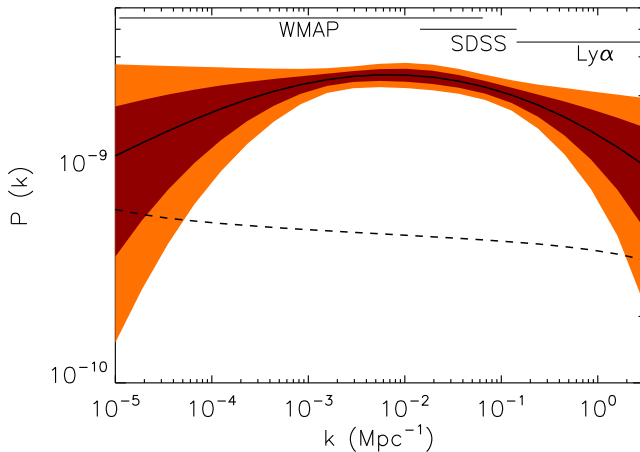
New results

Conclusion



WMAP3, from Spergel et al, astro-ph/0603449

SDSS-LRG5, from Percival et al, astro-ph/0608636

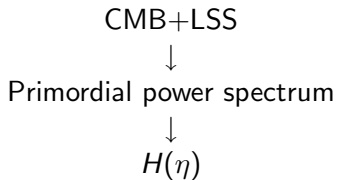


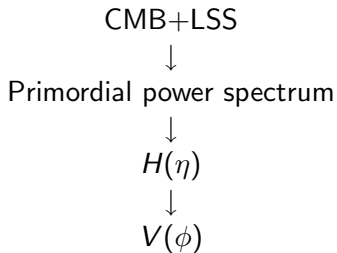
Taken from Easter & Peiris, astro-ph/0609003.

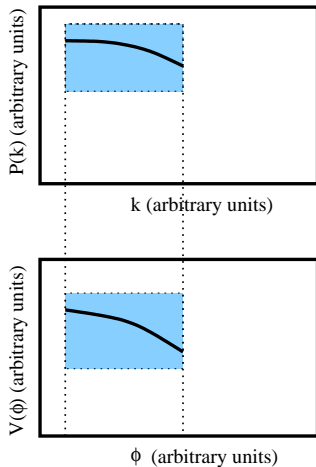
CMB+LSS

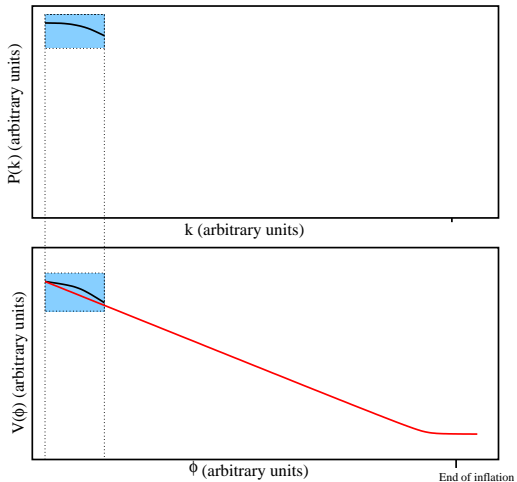


Primordial power spectrum

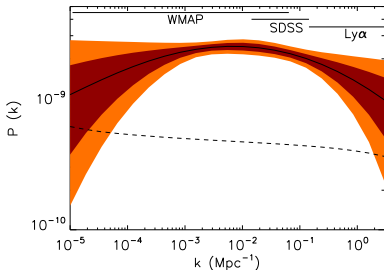




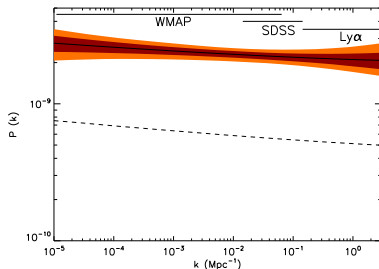




Approximations:

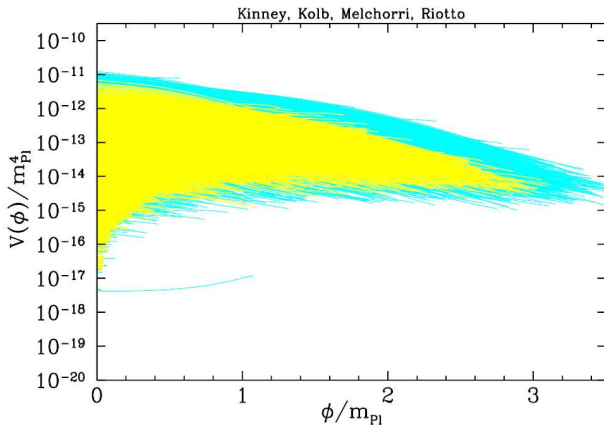


Fitting $P(k) = k^{(n_s-1+\dots)}$.



Fitting $P(k) = k^{(n_s-1+\dots)}$,
 selecting SR-inflationary models
 with $N > 30$. In this case
 WMAP1 already constrains V'''' .
 (Caprini et al. 2002)

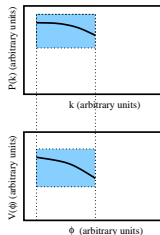
Taken from Easter & Peiris, astro-ph/0609003.



Taken from Kinney et al., astro-ph/0605338.

Directly fit the inflaton potential, numerically, using COSMOMC^I and our own freely available module^{II}.

$$\text{CBM} + \text{LSS} \\ \updownarrow \\ V(\phi)$$

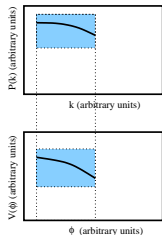


^ILewis & Bridle, 2002

^{II}see astro-ph/0703625

Directly fit the inflaton potential, numerically, using COSMOMC^I and our own freely available module^{II}.

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Result applies to any theory of inflation which, during the observable window, has effectively one scalar degree of freedom.

^ILewis & Bridle, 2002

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Directly fit the inflaton potential, numerically

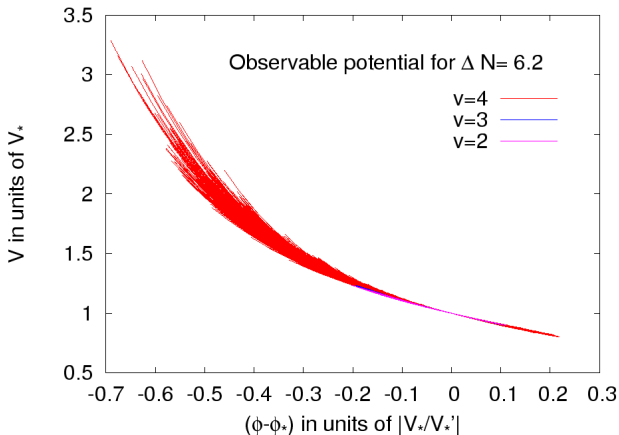
+ self-consistent
 tensor parameters:

$$n_T = -r/8,$$

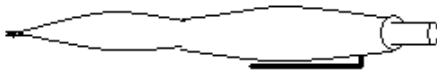
$$\alpha_T =$$

$$n_T [n_T - n_S + 1]$$

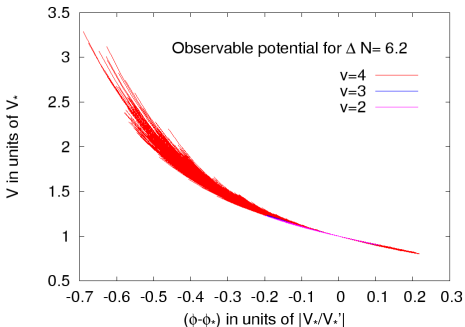
Slow Roll	Numerical potential fit
$\Omega_b h^2$	$\Omega_b h^2$
$\Omega_{cdm} h^2$	$\Omega_{cdm} h^2$
θ	θ
τ	τ
$\ln[10^{10} \mathcal{P}_{\mathcal{R}}^{k_*}]$	$\ln \left[\frac{128\pi 10^{10} V_*^3}{3V_*'^2 m_P^6} \right]$
r	$\left(\frac{V_*'}{V_*} \right)^2 m_P^2$
n_S	$\frac{V_*'''}{V_*} m_P^2$
α_S	$\frac{V_*'''}{V_*} \frac{V_*'}{V_*} m_P^4$
β_S	$\frac{V_*''''}{V_*} \left(\frac{V_*'}{V_*} \right)^2 m_P^6$



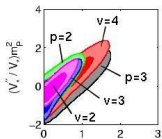
Did anyone see my pencil? (Parker..)



Conclusion



- ▶ The data allows departure from Slow Roll
- ▶ Previously obtained info on $V(\phi)$ depends on strong assumptions.
- ▶ Conservative analysis now gives $V(\phi)$ up to $V'''(\phi)$
- ▶ Hint to go to one order higher in SR



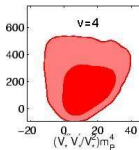
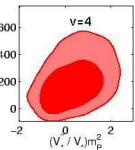
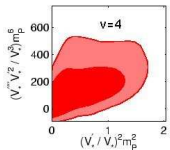
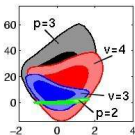
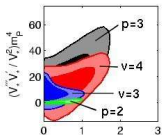
$p=2$ - A_S, n_S

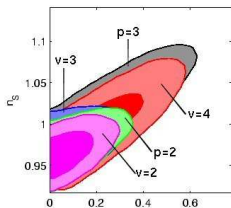
$p=3$ - A_S, n_S, α_S

$v=2$ - V'_*, V''_*

$v=3$ - V'_*, V''_*, V'''_*

$v=4$ - $V'_*, V''_*, V'''_*, V''''_*$





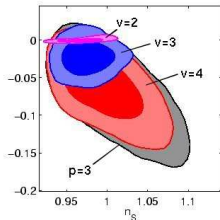
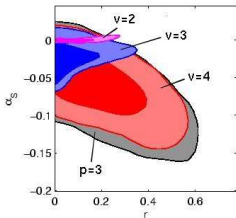
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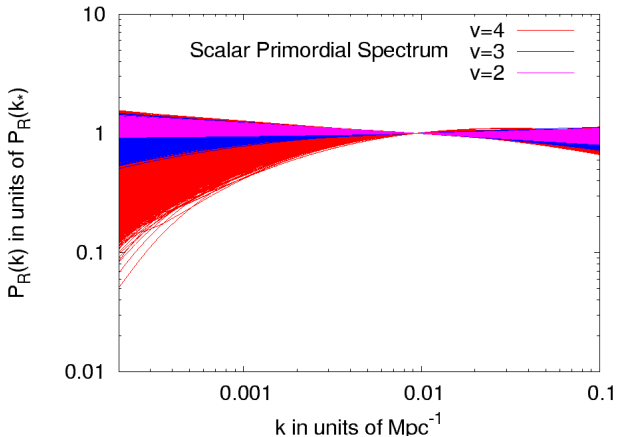
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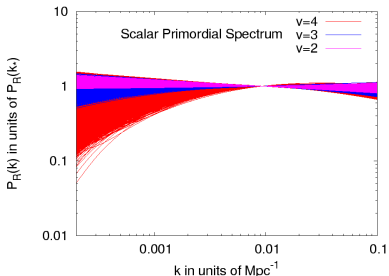
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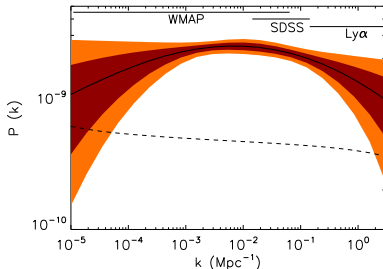




Approximations:

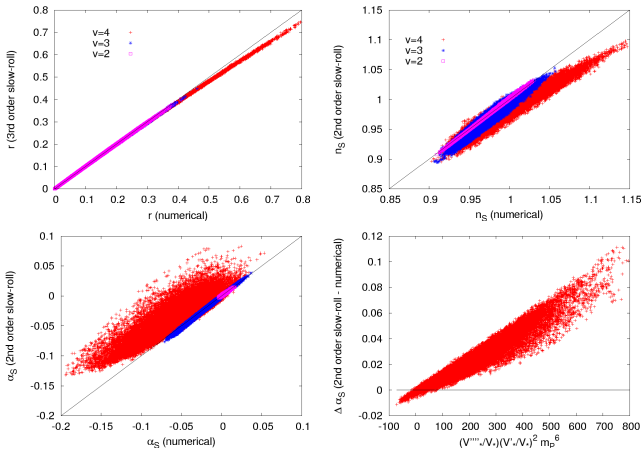


Fitting $V(\phi - \phi_*)$.



Fitting $P(k) = k^{(n_S - 1 + \dots)} \cdot a$

^aTaken from Easter & Peiris, astro-ph/0609003.



$$\dot{\phi} = -\frac{m_P^2}{4\pi} H'(\phi)$$
$$[H'(\phi)]^2 - \frac{12\pi}{m_P^2} H^2(\phi) = -\frac{32\pi^2}{m_P^4} V(\phi)$$

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$$\partial_\eta^2 \mu_{S,T} + \left[k^2 - \frac{\partial_\eta^2 z_{S,T}}{z_{S,T}} \right] \mu_{S,T} = 0$$

$$\mathcal{P}_{\mathcal{R}}(k) = \frac{k^3}{8\pi^2} \left| \frac{\mu_S}{z_S} \right|^2$$

$$\mathcal{P}_h(k) = \frac{2k^3}{\pi^2} \left| \frac{\mu_T}{z_T} \right|^2$$

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Primordial power spectrum of curvature perturbations:

$$\ln \frac{\mathcal{P}_{\mathcal{R}}(k)}{\mathcal{P}_{\mathcal{R}}(k_*)} = (n_S - 1) \ln \frac{k}{k_*} + \frac{\alpha_S}{2} \ln^2 \frac{k}{k_*} + \frac{\beta_S}{6} \ln^3 \frac{k}{k_*} \dots,$$

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Slow-roll approximation of inflation:

$$n_S - 1 = -2\epsilon_1 - \epsilon_2 - 2\epsilon_1^2 - (2C + 3)\epsilon_1\epsilon_2 - C\epsilon_2\epsilon_3,$$

$$\alpha_S = -2\epsilon_1\epsilon_2 - \epsilon_2\epsilon_3$$

$$\epsilon_0 \equiv H(N_i)/H(N)$$

$$\epsilon_{n+1} \equiv \frac{d \ln |\epsilon_n|}{dN}, \quad n \geq 0.$$

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Slow-roll approximation of inflation:

$$\epsilon_n = \epsilon_n(n_S, \alpha_S, \dots)$$

$$V = \frac{3m_{\text{Pl}}^2 H^2}{8\pi} \left(1 - \frac{\epsilon_1}{3}\right)$$

$$V' = -\frac{3m_{\text{Pl}} H^2}{(4\pi)^{1/2}} \epsilon_1^{1/2} \left(1 - \frac{\epsilon_1}{3} + \frac{\epsilon_2}{6}\right)$$

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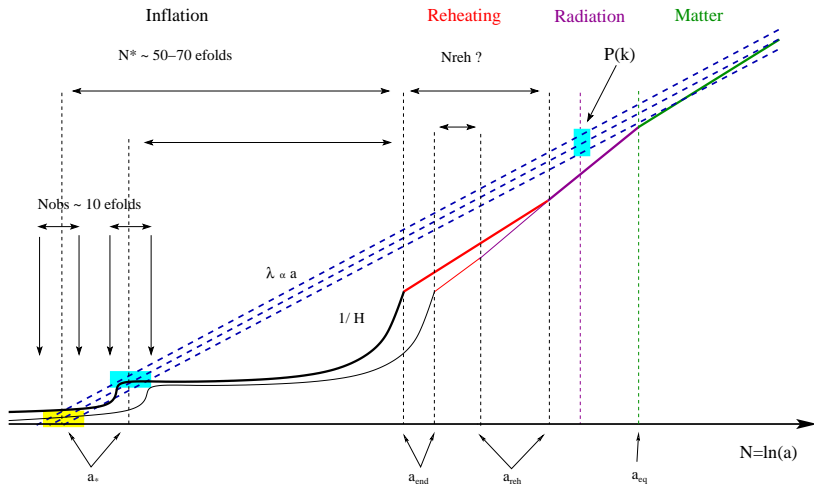
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 - ▶ Condition: $-d \ln H / d \ln a < 1$.
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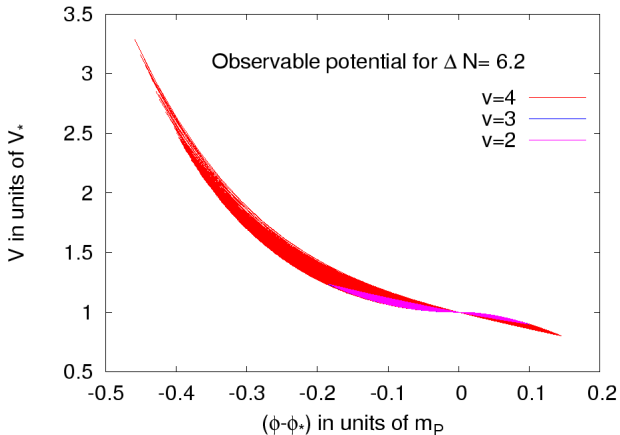
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- ▶ Fit to data using an MCMC.



¹¹taken from Ringeval, astro-ph/0703486



$$\lim_{k/aH \rightarrow \infty} \mu_{S,T}(\eta) = \frac{4\sqrt{\pi}}{m_{\text{Pl}}} \frac{e^{-ik(\eta-\eta_i)}}{\sqrt{2k}}$$

$$\mu_S(\eta) \equiv 2z_S \mathcal{R}$$

$$\mu_T(\eta) \equiv z_T h$$

$$z_S \equiv a \sqrt{2 - aa''/a'^2}$$

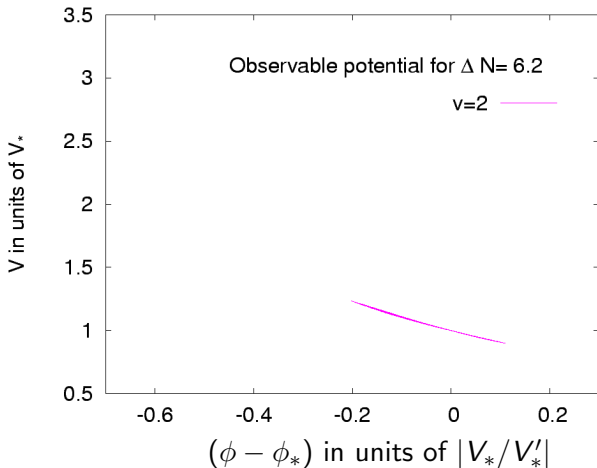
$$z_T = a$$

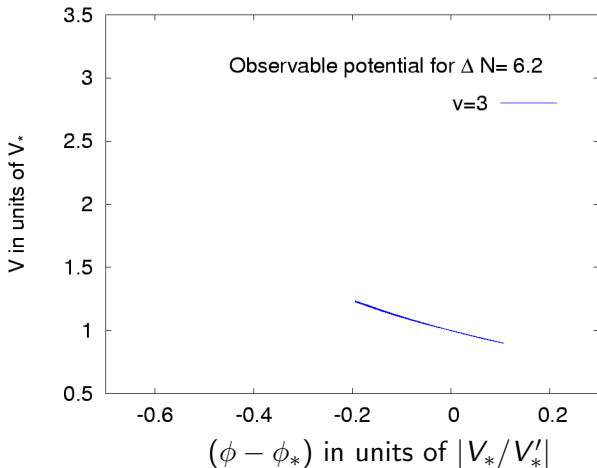
$$\frac{k_*}{a_I H_I} = \frac{k_*}{a_0 H_0}$$

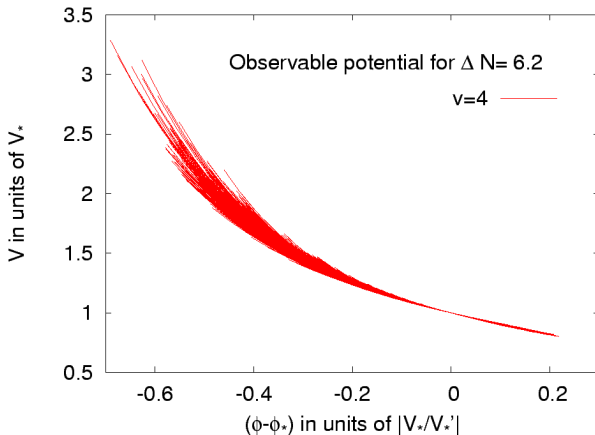
$$\frac{a_0}{a_I} = \frac{H_I}{H_0}$$

$$2N \equiv \ln \frac{a_0}{a_I} = \ln \frac{H_I}{H_0}$$

$$N \simeq 60 + \frac{1}{2} \ln \frac{H_I}{10^{13} \text{GeV}}$$







$$V = \frac{3m_{\text{Pl}}^2 H^2}{8\pi} \left(1 - \frac{\epsilon_1}{3}\right)$$

$$V' = -\frac{3m_{\text{Pl}} H^2}{(4\pi)^{1/2}} \epsilon_1^{1/2} \left(1 - \frac{\epsilon_1}{3} + \frac{\epsilon_2}{6}\right)$$

$$\frac{V''}{3H^2} = 2\epsilon_1 - \frac{\epsilon_2}{2} - \frac{2\epsilon_1^2}{3} + \frac{5\epsilon_1\epsilon_2}{6} - \frac{\epsilon_2^2}{12} - \frac{\epsilon_2\epsilon_3}{6}$$

$$V''' = \frac{12m_p^2 H^2 \sqrt{\pi}}{\sqrt{\epsilon_1}} \left(2\epsilon_1^2 - \frac{3\epsilon_2\epsilon_1}{2} + \frac{\epsilon_2\epsilon_3}{4}\right).$$