Selected topics in

extra-dimensional cosmology

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Plan

Overview of the different models.

Motivations

ADD

Randall-Sundrum models

Flat TeV X-dim

Dark Matter from X-dim: the different candidates

cosmology of flat TeV X-dim: SuperWIMPS, Radion cosmology

WIMP KK dark matter in UED

RS1 cosmology

radion stabilisation in RS. Goldberger-Wise mechanism

EW phase transition from radion stabilisation

properties, signatures and main exp. constraints for each of these models

> 2nd lecture (blackboard)

1st lecture (laptop) Many other subjects I am not covering

Inflation not discussed but covered by J. Cline's lectures and maybe by Lev Kofman

Brane cosmology (solutions to Einstein eqs, cosmological see for instance, hep-th/9905012 perturbations ... in D>4) (Binetruy, Deffayet, Langlois), or Kofman al: hep-ph/0404141

Self-tuning approach to the Cosm. Constant pb.

Asymmetric warped space-time (spatial and time scales have see for instance Csaki et al: different warp factors --> violation of Lorentz invariance) gr-qc/0105114

DGP model (Dvali-Gabadadaze-Porrati) presented in Gabadadze lectures hep-ph/0308112 hep-th/0005016 for a critical analysis see for instance: hep-th/0404159

Other attempts on modified gravity at large distances Gregory-Rubakov-Sibiryakov th/0002072 Dvali-Gabadadze-Shifman th/0202174

Transplanckian effects

String gas cosmology

Deconstruction

. . .

Some references. I- Set of lectures

TASI lectures on extra dimensions and branes. <u>Csaba Csaki</u> : hep-ph/0404096

effectives theories, ADD, symmetry breaking in flat Xdim via orbifolds (EW,susy,GUTs), mediation of susy breaking, warped pheno

TASI lectures on electroweak symmetry breaking from extra dimensions. <u>Csaba Csaki</u>, Jay Hubisz, <u>Patrick Meade</u> hep-ph/0510275

more detailed look at gauge theories in Xdim, higgsless models, fermions in Xdim, EW precision observables

TASI 2004 lectures on the phenomenology of <u>more phenomenological + Universal Extra Dimensions (UED)</u> extra dimensions. <u>Graham D. Kribs.</u> hep-ph/0605325

Les Houches lectures on warped models and holography. <u>Tony Gherghetta</u> hep-ph/0601213

Cargese Lectures on Extra Dimensions. <u>R. Rattazzi</u>: hep-ph/0607055

Large and infinite extra dimensions: An Introduction. <u>V.A. Rubakov</u> : hep-ph/0104152

ICTP lectures on large extra dimensions. Gregory Gabadadze : hep-ph/0308112

Tasi 2004 lectures: To the fifth dimension and back. Raman Sundrum hep-th/0508134

An Introduction to extra dimensions. <u>Abdel Perez-Lorenzana</u> hep-ph/0503177 warped models, susy warped, warped GUTs, AdS/CFT, holography

effective actions, ADD,RS,Goldberger-Wise stabilization, AdS/CFT, holography oduction. similar as above + localization of fermions and gauge fields + Cosm. Const. + modified gravity + lorentz violation

KK theories, ADD, warped models + DGP model

effective theories, orbifolds and chirality, radion stabilisation, cosm. constant pb, warpeds models

effective actions ... + neutrino mass models, split fermions, 6D models

II- Some original references

Phenomenology, astrophysics and cosmology of theories with submillimeter dimensions and TeV scale quantum gravity. <u>Nima Arkani-Hamed</u>, <u>Savas Dimopoulos</u>, <u>G.R. Dvali</u> hep-ph/9807344

An Alternative to compactification. Lisa Randall, Raman Sundrum hep-th/9906064

A Large mass hierarchy from a small extra dimension. Lisa Randall, Raman Sundrum hep-ph/9905221

Holography and phenomenology. Nima Arkani-Hamed, Massimo Porrati, Lisa Randall hep-th/0012148

Comments on the holographic picture of the Randall-Sundrum model. <u>R. Rattazzi</u>, <u>A. Zaffaroni</u> hep-th/0012248

II- On cosmology

On dark matter from extra dimensions:

section 25 of Les Houches physics at TeV colliders 2005 beyond the standard model working group: Summary report. hep-ph/0602198 (section 23 is on UED)

On RS1 cosmology:

Cosmology of brane models with radion stabilization. Csaba Csaki, Michael Graesser, Lisa Randall, John Terning. hep-ph/9911406

Holography and the electroweak phase transition. Paolo Creminelli, Alberto Nicolis, Riccardo Rattazzi hep-th/0107141

Order rho² corrections to Randall-Sundrum I cosmology. James M. Cline, Jeremie Vinet hep-th/0201041

also relevant:

Exact identification of the radion and its coupling to the observable sector. Lev Kofman, Johannes Martin, Marco Peloso hep-ph/0401189

Why consider theories with extra dimensions?

- × D=3+1 is not a prediction in Einstein's theory
- × Only string theory predicts the number of dimensions i.e D=1+9(10).

We have to 'hide' extra dimensions Indeed, $F \sim \frac{1}{r^2}$: only in 3 dimensions!

X Easy to hide extra dimensions if they are compact and tiny:

$$F \sim \frac{1}{r^{2+n}} \text{ for } r < r_c$$
$$F \sim \frac{1}{r^2} \text{ for } r \ge r_c$$

X Not only are extra dimensions allowed but they could also be useful to help us resolve the big puzzles of 4D physics...

A few puzzles in Beyond the Standard Model Physics

electroweak sector

pbs related to fermions

Grand Unification?

cosmological pbs

-Neutrino masses and hierarchy in fermion masses -Proton stability $\frac{m_{\nu}}{m_{\nu}} \sim 10$ $m_{\nu} \sim 10^{-14}$ -Flavour problem (FCNCs)

-Hierarchy problem $\frac{M_{\rm EW}}{M_{\rm Pl}} \sim 10^{-16}$ -Symetry breaking (electroweak symmetry, supersymmetry)

-Unification of couplings

-Cosmological constant -Inflation -Dark energy, duintessence $\rho_{\Lambda} \simeq (\text{meV})^4$ -Dark matter $\Omega_{\text{mat noire}} \simeq 25\%$ -Baryogenesis $\frac{n_b}{n_\gamma} \sim 10^{-10}$

Physics beyond the Standard Model implies new

degrees of freedom.

Why not extra dimensions ?

Which size for extra

Pimensions?

Experimental constraints:

Roughly:

- If Standard Model fields propagate in extra dimensions $\Rightarrow R < (TeV)^{-1} \sim 10^{-19} m$

-If only gravity propagates in extra dimensions \Rightarrow R < mm

If we assume that ALL fields propagate in ALL extra dimensions (which moreover have ALL the same size):

 $\begin{array}{c} M_{\rm Pl}^2 \sim R^n M_*^{n+2} \\ \frac{1}{g_4^2} \sim R^n M_*^n \end{array} \end{array} \begin{array}{c} R \sim \frac{g_4^{\frac{n+2}{n}}}{M_{\rm Pl}} & \mbox{i.e.} & R \sim M_{\rm Pl}^{-1} \end{array} \end{array}$

Things change if, in particular, fields are LOCALIZED in extra dimensions.

and this brings us to the ADD idea: (Arkhani-Hamed, Dimopoulos, Dvali '98) The Planck scale is no longer a fondamental scale but an effective scale:

If M_* is the fundamental scale and $M_* \sim {
m TeV}$

then $n = 2 \Rightarrow R^{-1} \sim \text{meV} \sim \text{mm}^{-1}$



- "Large" flat extra dimensions(can be almost macroscopic in size) where only gravity propagates.

- The Standard Model is localized in 3 dimensions.

 $R \sim meV^{-1} \sim mm$

: no stark disagreement with experiments and observations.



Gravity appears weak because it is 'diluted' in extra dimensions

Extra-dimensional gravitons look to us as a "tower" of massive gravitons with masses regularly-spaced in n/R



Signatures at colliders Observing quantum gravity at the LHC

Each graviton taken individually has a coupling suppressed in 1/Mpl and the production of a single graviton est totally negligible.

However, the cross section to produce a collection of massive gravitons is amplified due to the vary large number of gravitons.

 Δ m~meV -> continuum of states $\sigma \sim rac{(ER)}{M^2} \sim rac{E}{M^{n+2}}$

The effective scale suppressing the coupling is in fact M_{st} and not $M_{
m Pl}$



Signatures at colliders Observing quantum gravity at the LHC ⇒ Direct production of gravitons

gravitons escape from the brane ->invisible KK graviton signature is monojet + missing energy continuum of states: mass distribution is a continuun

>>virtual graviton exchange

new reaction $g \ g \rightarrow l^+ l^-$ + deviations with respect to standard processes (interference with the amplitude in the SM)



black hole production

For $\sqrt{s}>>M_*$ semi classical description becomes adequate as $r_s>>M_*^{-1}$

quantum-gravity effects are subleading with respect to classical gravitational effects

When the impact parameter $\, b < r_s \,$ we expect black hole formation

The LHC : a black hole factory

Black hole production of mass $M_{\rm BH} = \sqrt{s}$ if the impact parameter of the collision is smaller than the Schwarzschild radius.

 r_H

i.e geometrical approximation $\sigma \sim \pi r_H^2$ (transplanckian energies) $\sqrt{s} > M_*$

Schwarzschild radius in 4+n dimensions: 4+n generalization of $ds^2 = (1 - \frac{GM}{r})dt^2 - \frac{dr^2}{1 - GM/r} + r^2d^2\Omega$

$$\sim \left(\frac{M_{\rm BH}}{M_*}\right)^{\frac{1}{n+1}} \frac{1}{M_*}$$

 $\sigma \sim \pi r_H^2 \sim \frac{1}{M_*^2} \left(\frac{M_{BH}}{M_*}\right)^{\frac{2}{n+1}} \sim \text{TeV}^{-2}$

Evaporation via Hawking radiation: $T_{H} = \frac{n+1}{4\pi r_{H}} \in [80-600] \text{ GeV for n=1...7}$ $\tau \sim \frac{1}{M_{*}} \left(\frac{M_{BH}}{M_{*}}\right)^{\frac{n+3}{n+1}} \sim 10^{-26} \text{sec}$



The mass gap is 1/R.

For n=2 and M_* =1 TeV, this is meV.

BBN energy is MeV and the number of KK gravitons which are kinematically accessible is more than 10^18!

problem: Too much energy is released into KK gravitons.

At $T \gtrsim 1/R$ the number of KK modes which are kinematically accessible is $(TR)^n$ The cross section for graviton production from brane thermal processes is $\sigma \sim \frac{(TR)^n}{M_{\rm Pl}^2} \sim \frac{T^n}{M_*^{n+2}}$ $\Gamma_G = \langle n_\gamma \sigma_{\gamma\gamma \to G} v \rangle \sim \frac{T^{n+3}}{M_*^{n+2}}$ no backward processes: $\frac{dn_G}{dt} \sim n_\gamma \Gamma_G \rightarrow n_G \sim n_\gamma \Gamma_G H^{-1} \sim \frac{T^{n+4}M_{Pl}}{M^{n+2}}$ $\frac{n_G}{n_\gamma} < 1 \rightarrow T_* < \left(\frac{M^{n+2}}{M_{Pl}}\right)^{1/(n+1)}$ cooling bound

which can also be derived by demanding that the cooling of the universe due to evaporation of KK gravitons in extra dimensions be smaller than the cooling due to expansion:

$$\frac{d\rho}{dt}|_{evap} \sim -\frac{T^{n+7}}{M_*^{n+2}} < \frac{d\rho}{dt}|_{exp} \sim -3H\rho \sim -3\frac{T^6}{M_{Pl}}$$

BBN bound

Once KK gravitons are produced, they behave as matter of mass T (the probability to interact with the thermal bath on the brane is very small) and their energy density $\rho_G \sim T \times n_G$ redshifts as $1/R^3$: $\rho_G(T = MeV) \sim \rho_G(T) \left(\frac{MeV}{T}\right)^3$

slightly stronger than the overcooling bound

The two previous bounds apply to 4D particles with 1/TeV coupling. Moreover, the specificity of our gravitons is that the probability that they interact with the SM wall is very tiny: the energy stored in them can easily overclose the universe.

 $\implies \frac{\rho_G}{\rho_\gamma}|_{BBN} \sim \frac{T}{1 \ MeV} \frac{\rho_G}{\rho_\gamma}|_T \sim \frac{T}{1 \ MeV} \frac{T_*^{n+1}M_{Pl}}{M^{n+2}} \implies T < \left(\frac{10^{-3}M^{n+2}}{M_{Pl}}\right)^{1/(n+2)}$

Overclosure bound

 T^3

: Decay of a single graviton is indeed suppressed by 1/Mpl

 $ho_G \sim r$

$$\Gamma \sim \left(\Gamma_{\text{near wall}} \sim \frac{T^{n+3}}{M_*^{n+2}}\right) \times \begin{array}{c} \text{proba to be} \\ \text{near wall} \sim - \\ \sim \left(\frac{T^{-1}}{R}\right)^n \\ \sim \left(\frac{T^{-1}}{R}\right)^n \\ \text{wavelength ~1/T} \end{array}$$

The energy density stored in KK gravitons produced at temperature T, redshifts as 1/R^3 so $\frac{\rho_G}{T^3} \sim \frac{T^{n+2}M_{Pl}}{M^{n+2}} = \text{constant}$ and we require that

$$\frac{\rho_G}{T^3} < \frac{\rho_c}{T_0^3} \quad \text{where} \quad \frac{\rho_c}{T_0^3} \sim 3 \times 10^{-9} \text{GeV}$$
$$\implies \qquad T < \left(\frac{10^{-21} M^{n+2}}{M_{Pl}}\right)^{1/(n+2)}$$

bound from diffuse photon background

The fraction of KK gravitons produced at temperature T, with lifetime $au(T) \sim 10^{10} \mathrm{yr} \times \left(\frac{100 \mathrm{MeV}}{T}\right)^3$ which have already decayed is $\left(\frac{100 \mathrm{MeV}}{T}\right)^3$

The resulting number density of photons is



Constraint from COMPTEL data leads to $T < (10^{-40} GeV^3 M^{n+2})^{1/(n+5)}$

(Gamma ray observations in the MeV range)

Summary of constraints



Conclusion: Difficulty to implement leptogenesis/baryogenesis in this context. Cut off is TeV: How to make inflation natural?

Very strong constraints from astrophysics Ocosmology

X Cooling of supernovae and red giants due to graviton emission.

 $(M_* \ge 30 \text{ TeV} (n=2))$

X Distorsion in CMB due to graviton decay (primary or secondary)

X heating of neutron stars due to KK graviton decay

X Overclosure of the universe by gravitons

X Reheating temperature of the universe has to be very low otherwise gravitons evaporate into the bulk

 $T_{RH} \le T_* \qquad T_* \sim \left(\frac{M_*^{n+2}}{M_{\rm Pl}}\right)^{\frac{1}{n+1}} \qquad M_* = 1 \text{ TeV} \implies \begin{array}{l} n=2 \ \text{->T} \sim 0.7 \text{GeV} \\ n=6 \ \text{->T} \sim 317 \text{ GeV} \\ m=4 \ \text{->T} \sim 10 \text{ MeV} \\ n=4 \ \text{->T} \sim 1 \text{ GeV} \\ n=6 \ \text{->T} \sim 7 \text{ GeV} \end{array}$

Note: These constraints are relaxed if compact extra dimensions are hyperbolic rather than toroidal Kaloper et al hep-ph/0002001

 $(M_* \ge 110 \text{ TeV} (n=2))$

 $(M_* \ge 1700 \text{ TeV} (n=2))$

 $M_* \ge 8 \text{ TeV}$

problems related to the radion



The energy stored in radion oscillations overclose the universe (similar to axions)...

The Randall-Sundrum model A complete solution to the hierarchy problem

Non flat geometry but Anti de Sitter (non factorisable geometry: "warping")

Fondamental scale : $M_{\rm Pl}$ $(k \sim r^{-1} \sim \Lambda_5 \sim M_{\rm Pl})$ (appearing in the 5D

effective action)

AdS space: the energy scale varies with position along 5th dimension



$$\Rightarrow M_{\rm EW} \sim M_{\rm Pl} e^{-k\pi r}$$

Natural stabilisation of radius (à la Goldberger-Wise) :

$$kr = \frac{4}{\pi} \frac{k^2}{m^2} \ln\left[\frac{v_h}{v_v}\right] \sim 10$$

Particle physics model building in warped space

2006 favourite set-up:

hierarchy pb
fermion masses
High scale unification
FRW cosmology



C = 1/2

 πR

Note: No susy here and many different realizations

Randall-Sundrum KK gravitons are very different from ADD

X Discrete spectrum with KK states non regularly spaced (proportional to the zeros of Bessel functions)

X $\Delta m \sim O(TeV)$ compared to $\Delta m \sim O(meV)$ in ADD

Remark: $r^{-1} \sim M_{\rm Pl}$ but $M_{KK} \sim {\rm TeV}$

X Each KK graviton couples as 1/TeV and not 1/ $M_{
m Pl}$





(The radion of Randall-Sundrum is also very different from that of ADD

x m ~ O(100) GeV x strongly coupled radion (in 1/TeV and not 1/ $M_{
m Pl}$)

no cosmological pb associated with the radion

the coupling of the radion to matter is similar to the coupling of Higgs to matter \Rightarrow radion phenomenology = Higgs phenomenology

... and possibility of Higgs-Radion mixing modifying Higgs phenomenology at LHC

Induced operator on the $\int d^4x \sqrt{g} \mathcal{R} \xi H^{\dagger} H \implies \mathcal{L} \supset 6\xi \gamma h \Box r$ TeV brane: radion-higgs kinetic mixing

After diagonalisation, modification of the Higgs standard couplings

3rd class of models: (Flat) extra dimensions at the TeV Solution to the hierarchy problem?

Same status as in SUSY : The higgs mass is stabilized against radiative corrections $(m_h \sim R^{-1})$ but it remains to explain why R^{-1} -TeV

same as in supersymmetry where we have to explain why $\,M_{
m SUSY}\,$ ~TeV

only bosons in bulk

<mark>R ~ Te</mark>V⁻¹ (flat)

2 sub-classes

of models

All SM fields in bulk "Universal" Extra Dimensions (UED) An important feature of the SM: Its fermions are chiral, i.e. the left and righthanded components of any Dirac fermion have different gauge quantum numbers.

While 5D fermions are 4 component-spinors.

This imposes constraints on the compactification of extra dimensions

The simplest compactification on a circle or a torus leads to non-chiral "vector-like" fermions.

The chirality of the 4D fermions has to be introduced by the boundary conditions at the end points of the interval.

A little detour : orbifold projections

-If D=5, the gauge field contains a quadri-vector V_{μ} and a scalar V_5 -5D fermions are 4 component- spinors and lead to mirror fermions at low energy. In order to eliminate unwanted degrees of freedom and get a chiral theory in 4D, we apply an orbifold projection

$$0 \underbrace{\left(\begin{array}{c} +\pi R \\ -\pi R \end{array}\right)}_{-\pi R} \Rightarrow 0 \underbrace{\left(\begin{array}{c} -\pi R \end{array}\right)}_{-\pi R} = 0 \underbrace{\left(\begin{array}{c} -\pi R \end{array}\right)}_{-\pi R}$$

We fold the circle (= identify y and -y) \Rightarrow we get a segment (0 and πR are the 2 fixed points)

The orbifold projection consists in imposing the following conditions

$$\begin{array}{ccc} y \to -y & A_{\mu} \to A_{\mu} & \Psi_L \to \Psi_L \\ & A_5 \to -A_5 & \Psi_R \to -\Psi_R \end{array}$$

At y= 0, πR : Even fields have Neumann boundary conditions $\partial_y \Phi = 0$ Odd fields have Dirichlet boundary conditions $\Phi = 0$ The zero mode of the odd field is projected out by the Z_2

Kaluza-Klein decomposition

Fourier expansion

for compactification

on a circle

$$\Phi(x^{\mu}, y) = \sum_{n} f^{n}(y) \Phi^{n}(x^{\mu})$$

Fields which are even (A_{μ}, Ψ_{L}) have a zero mode

$$\Phi(x^{\mu}, y) = \sqrt{\frac{1}{\pi R}} \Phi^0(x^{\mu}) + \sum_{n \ge 1} \sqrt{\frac{2}{\pi R}} \cos\left(\frac{ny}{R}\right) \Phi^n(x^{\mu})$$



Fields which are odd (A_5, Ψ_R) do not have a zero mode

$$\Phi(x^{\mu}, y) = \sum_{n \ge 1} \sqrt{\frac{2}{\pi R}} \sin\left(\frac{ny}{R}\right) \Phi^{n}(x^{\mu})$$

The "zero" mode fermions are chiral (and identified with the SM fermions). However, the other KK fermions are vector-like.

(a bit similar to supersymmetry where each SM particle is accompanied by partners)

Each left-handed (right-handed) SM fermion possesses a distinct Kaluza-Klein tower

Orbifold projections are intensively used to break symmetries:

One imposes different boundary conditions for the different components of a given multiplet

*** Electroweak symetry *** Grand unification symmetry *** Supersymmetry

`Universal' Extra Dimensions

Assumption: All SM propagate in extra dimension(s).

Translation Invariance along the 5th dimension \Rightarrow Conservation of de Kaluza-Klein number in interactions of the 4D effective theory.



Consequence: n=1 KK excitations can only be pair-produced

⇒Collider constraints are weak





Phenomenology very similar to supersymmetry with conserved R-parity

Every KK particle eventually decays into the LKP

This symmetry is broken by the <u>orbifold but there remains</u> a discrete symmetry called Kaluza-Klein parity : (-1)ⁿ

For instance:

⇒Odd-n KK modes can only couple by pairs ⇒The lightest KK mode (LKP) is stable

The Kaluza-Klein photon: an excellent candidate for dark matter

comparison between the *z* models

Fondamental scale of gravity



Kaluza-Klein Mass



 $M_* \sim \text{TeV}$

 $R^{-1} \sim \mathrm{meV}$

 $M_{KK} \sim R^{-1}$

RS (AdS)

 $M_* \sim M_{\rm Pl}$

 $R^{-1} \sim M_{\rm Pl}$

 $M_{KK} \sim TeV$

 $\begin{array}{c} \text{TeV} \\ \text{\& UED} \end{array} n = 1 \to M_* \sim 10^{13} \text{ GeV} \\ n = 2 \to M_* \sim 10^{10} \text{ GeV} \end{array} \quad R^{-1} \sim \text{TeV} \qquad M_{KK} \sim R^{-1} \end{array}$

Back to our to-do list ...

electroweak sector

 $\checkmark - \text{Hierarchy problem} \qquad \frac{M_{\rm EW}}{M_{\rm Pl}} \sim 10^{-16}$ $\checkmark - \text{Symmetry breaking (electroweak symmetry, supersymmetry)}$

pbs related to fermions \checkmark - Neutrino masses and hierarchy in fermion masses $\frac{m_{\nu}}{m_{\rm top}} \sim 10$ \checkmark - Flavour problem (FCNCs) $m_{\nu} \sim 10^{-14}$

Grand Unification? -Unification of couplings -Cosmological constant

-Baryogenesis $\frac{n_b}{n_\gamma} \sim 10^{-10}$

-Inflation

