

Resonant Polarimetry:

a way to non-invasive and fast beam polarimetry?!

Wolfgang Hillert

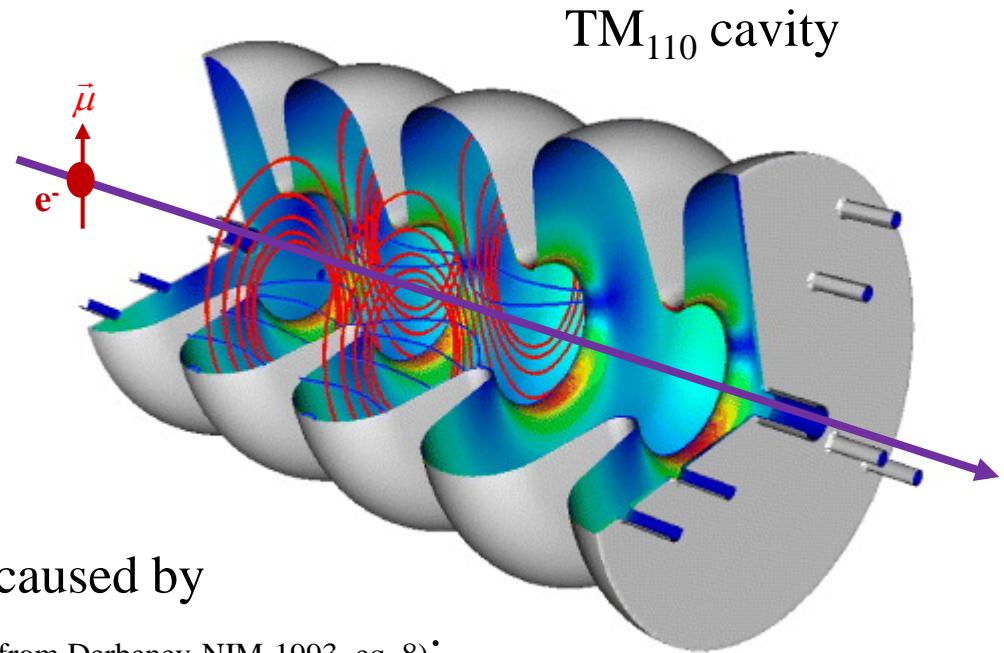
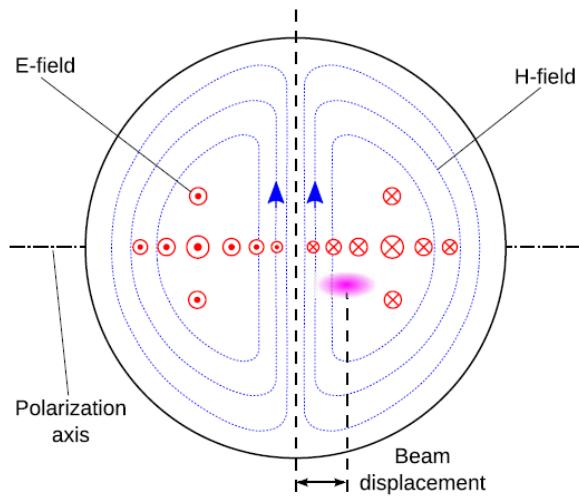


Physics Institute of Bonn University

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- 1. Functional Principle
 - 2. Relativistic Stern-Gerlach Force
 - 3. Cavity Modes
 - 4. Energy Transfer per Particle Passage
 - 5. Signal Power
 - 6. Example: Respol with TE_{011} , TE_{111}

Resonant Polarimetry

Principle Idea (Derbenev 1993):

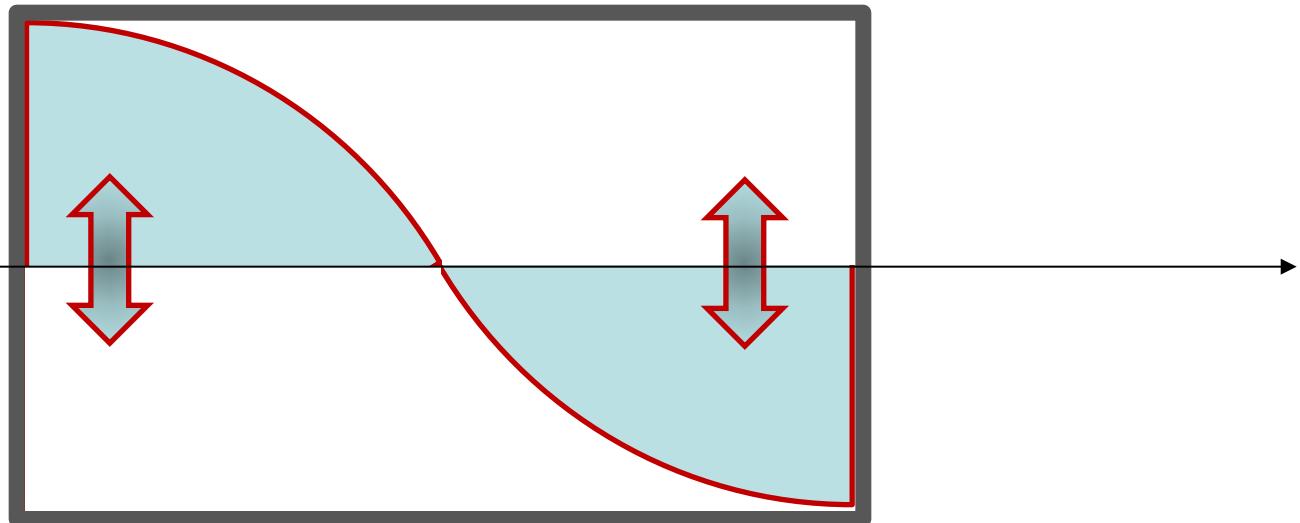


Coupling of the magnetic moment (caused by the spin) to the cavity's B-field (taken from Derbenev-NIM-1993, eq. 8):

$$W_c = \omega_c |a|^2 = \omega_c N^2 \left\langle \frac{e}{2mc\sqrt{2\omega_c}} \left(\left(G + \frac{1}{\gamma} \right) \vec{B}_\perp^c + \frac{1+G}{\gamma} \vec{B}_\parallel^c \right) \vec{e} \cdot e^{ik\theta} \right\rangle^2 \frac{\hbar^2 t^2}{4} P_e \sin^2 \alpha$$

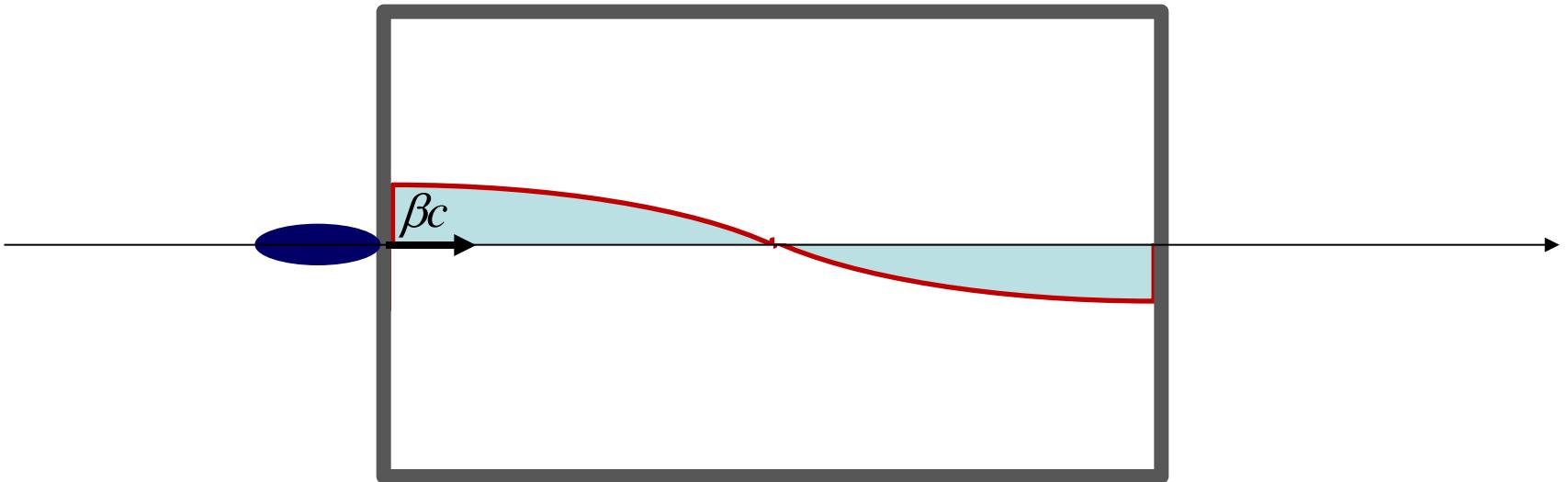
?Physical understanding? ?γ and G scaling?

Transverse Mode



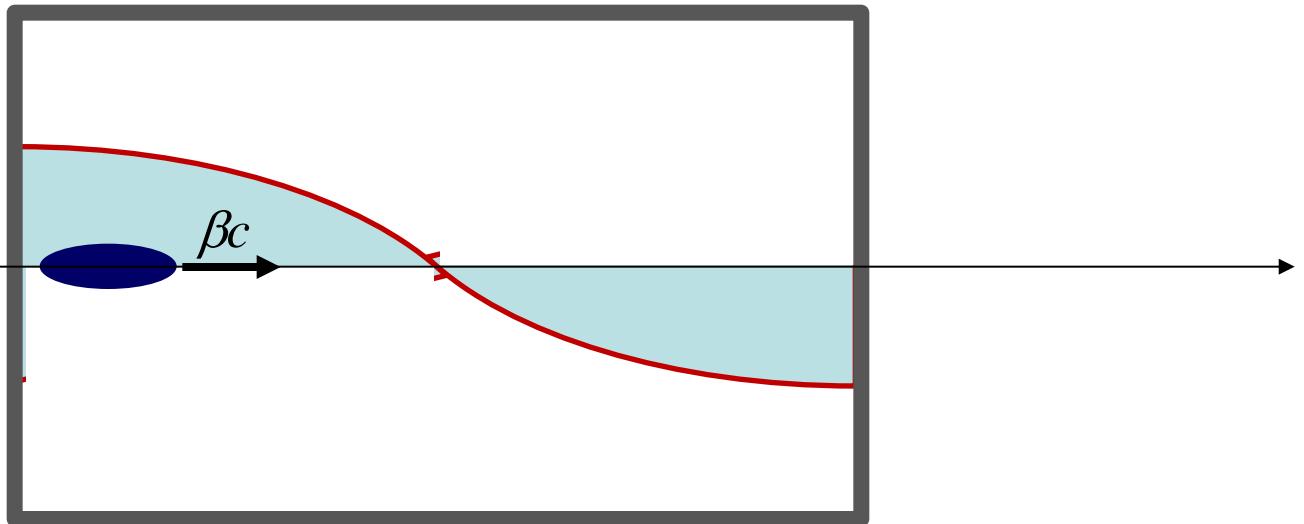
$$\Delta W = \int \frac{\partial}{\partial z} (\vec{\mu} \cdot \vec{B}) \cdot dz$$

Transverse Mode



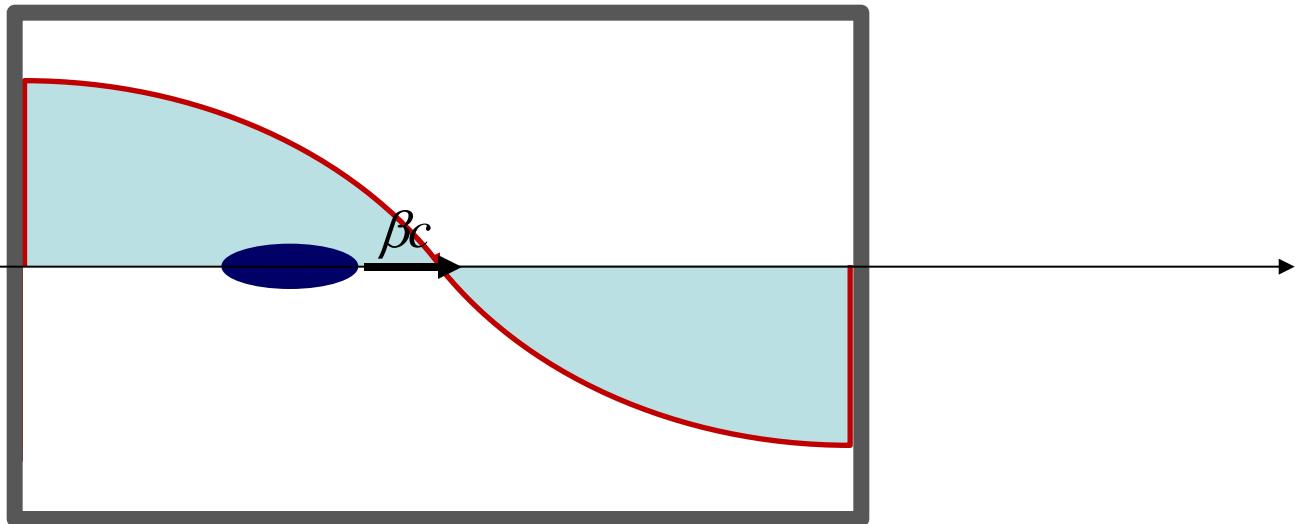
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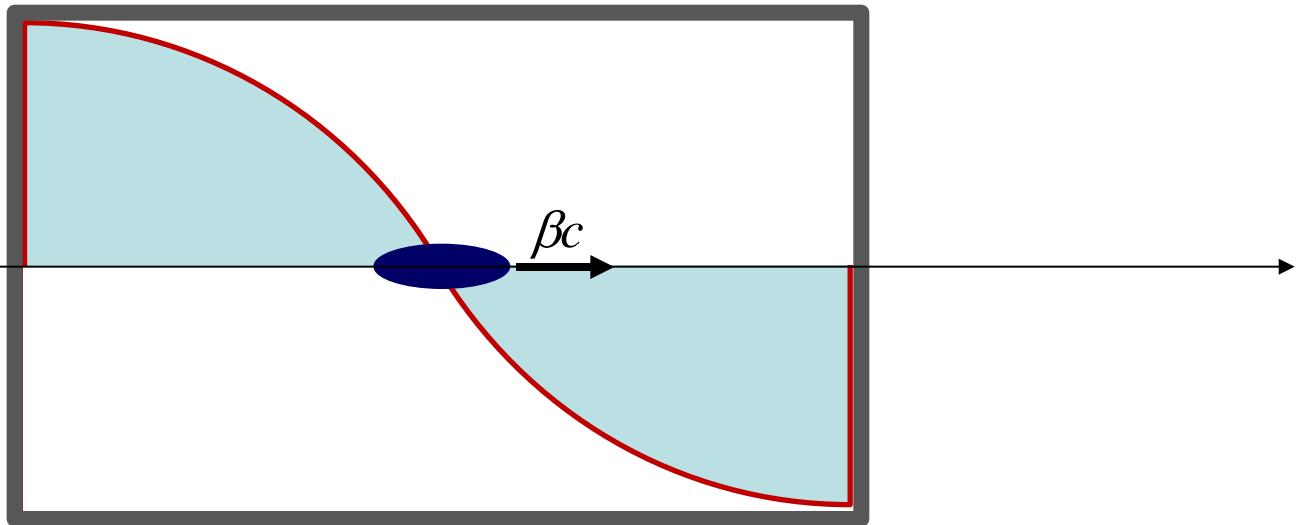
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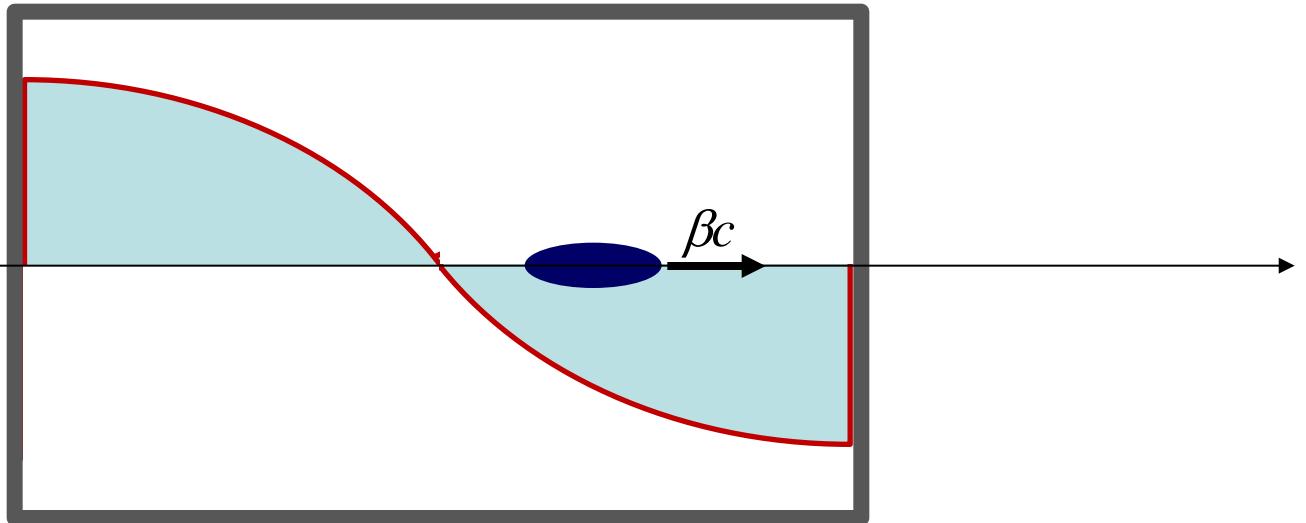
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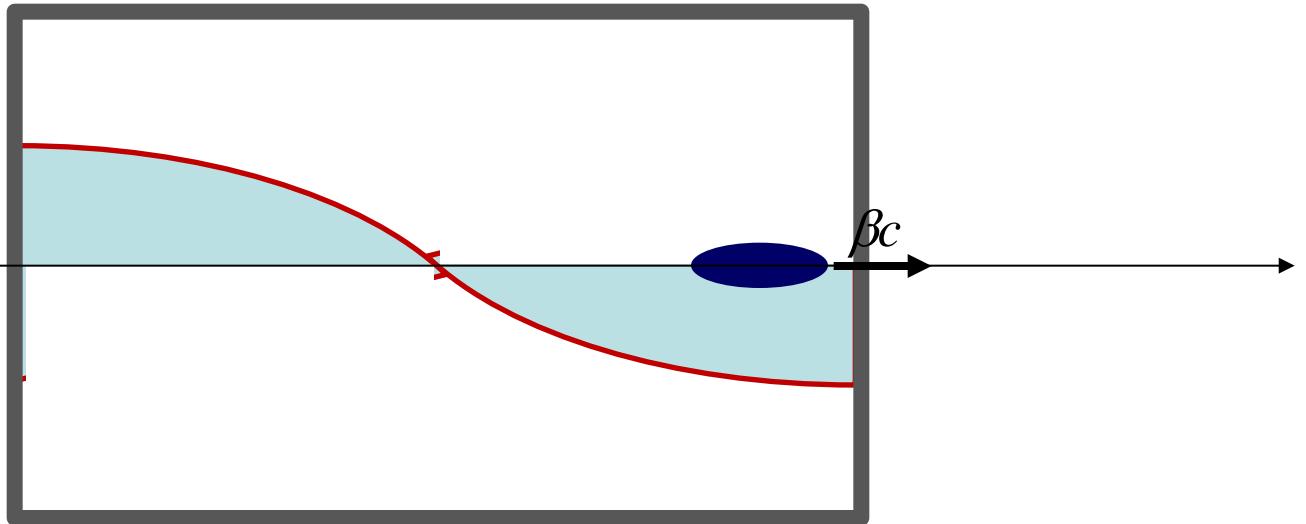
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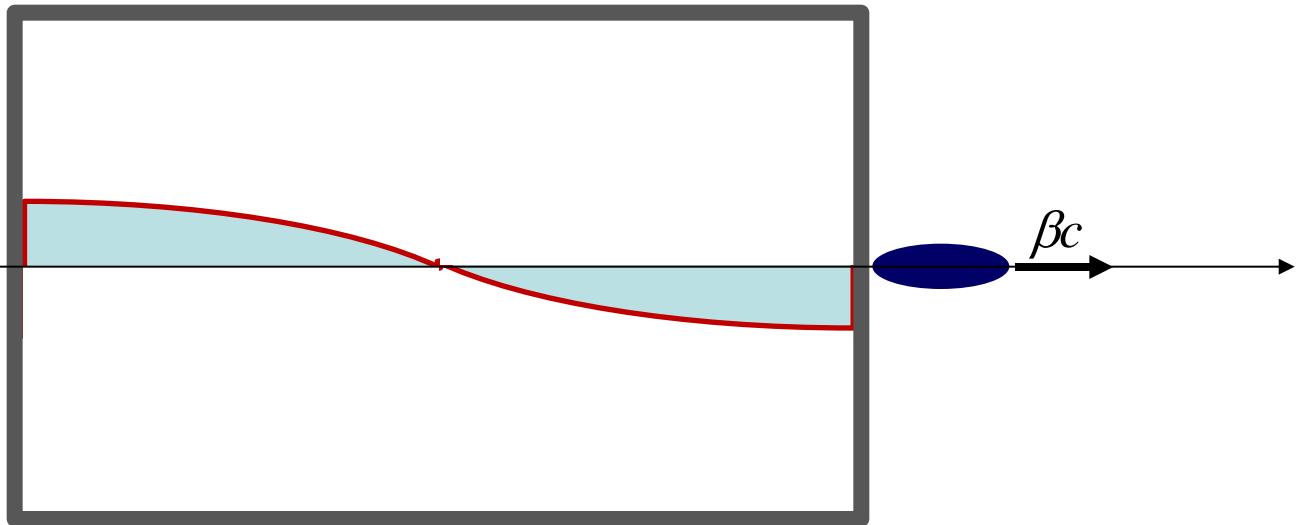
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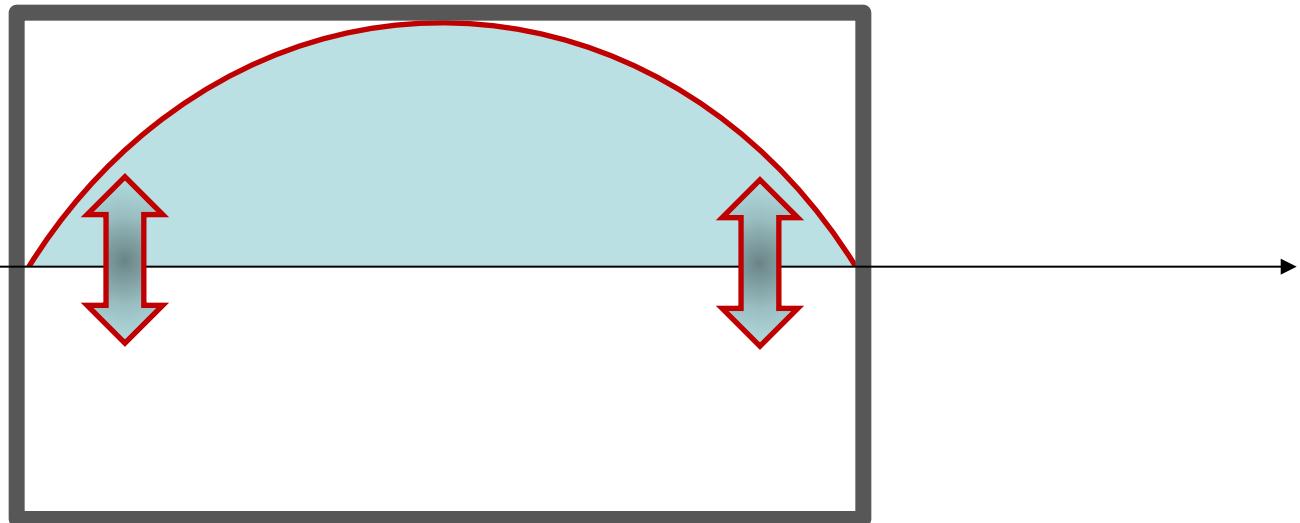
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Transverse Mode



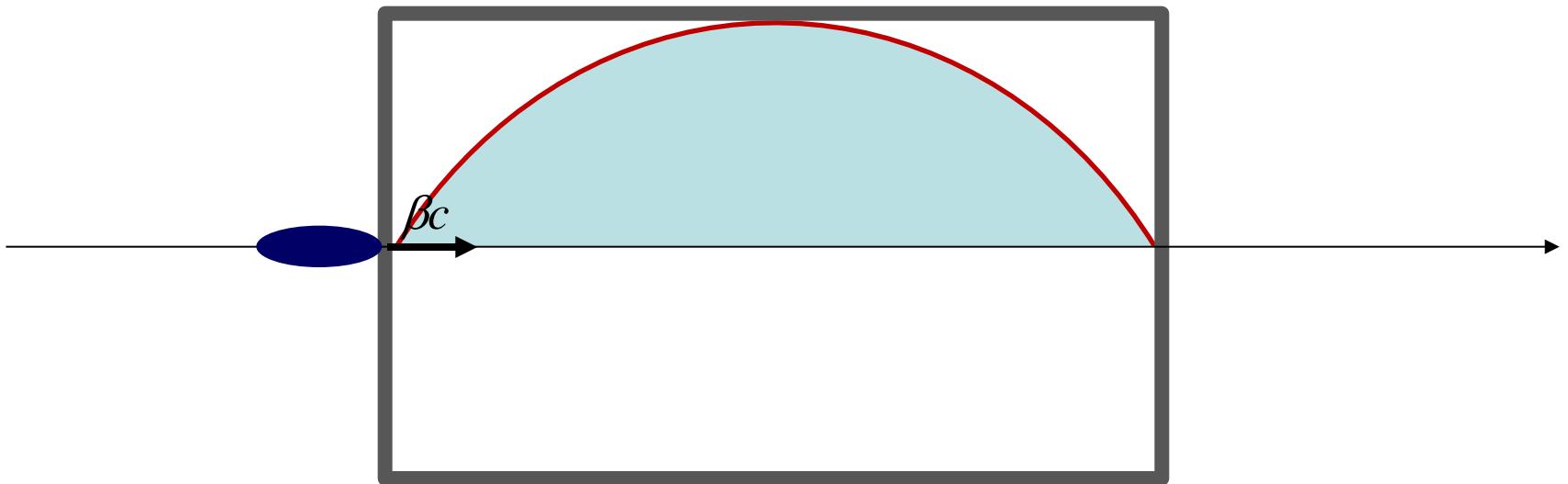
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Longitudinal Mode



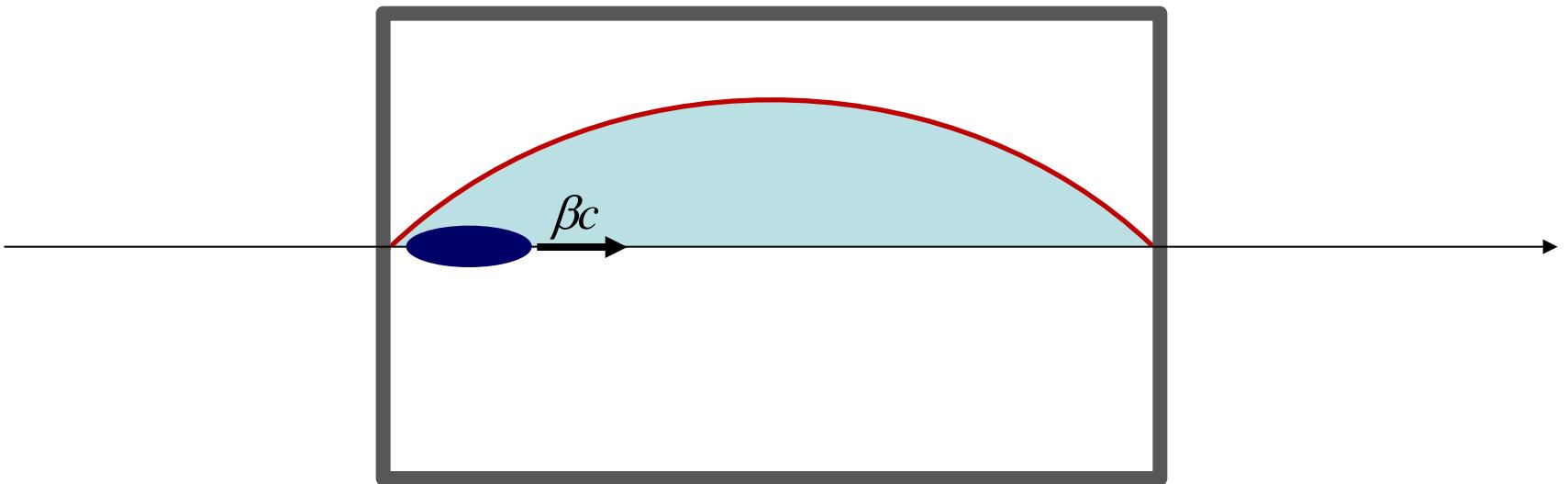
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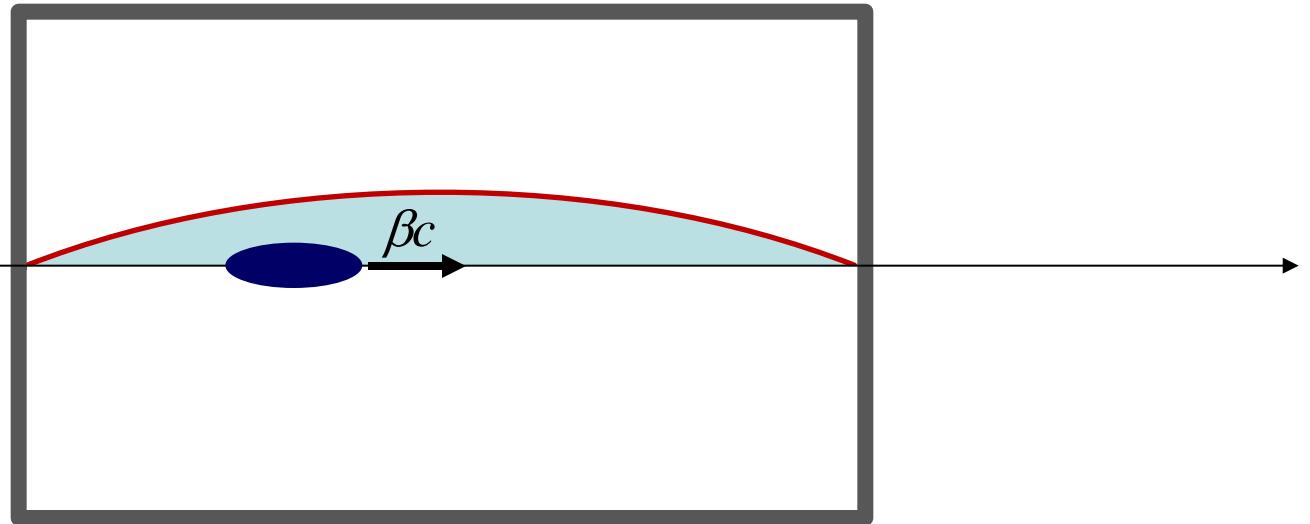
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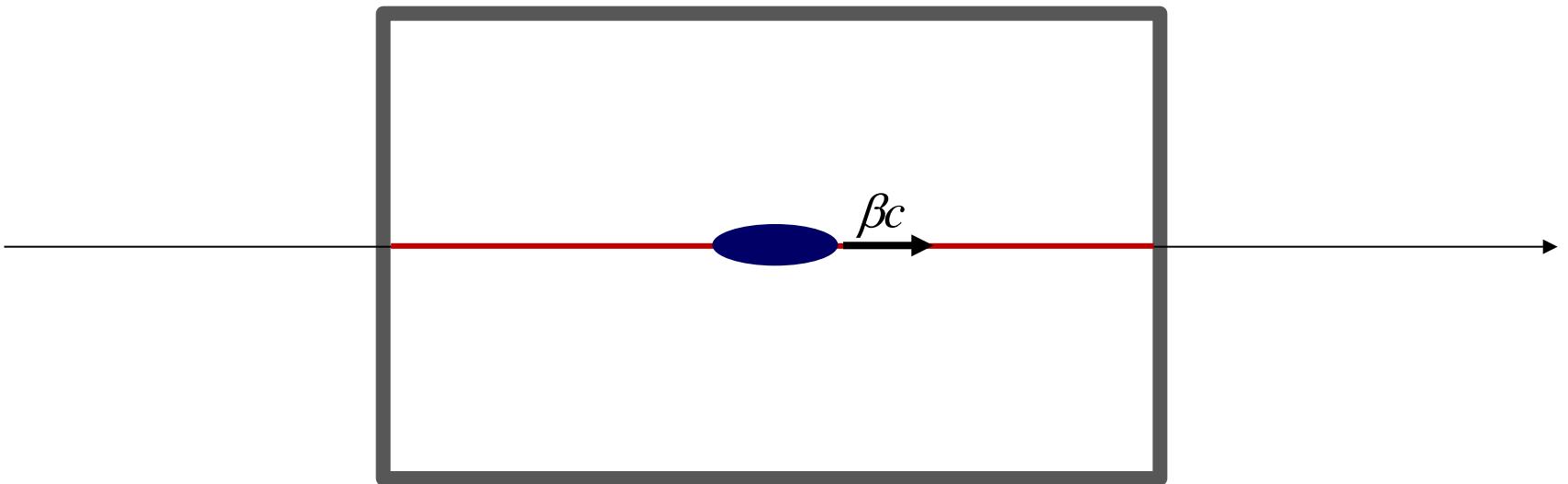
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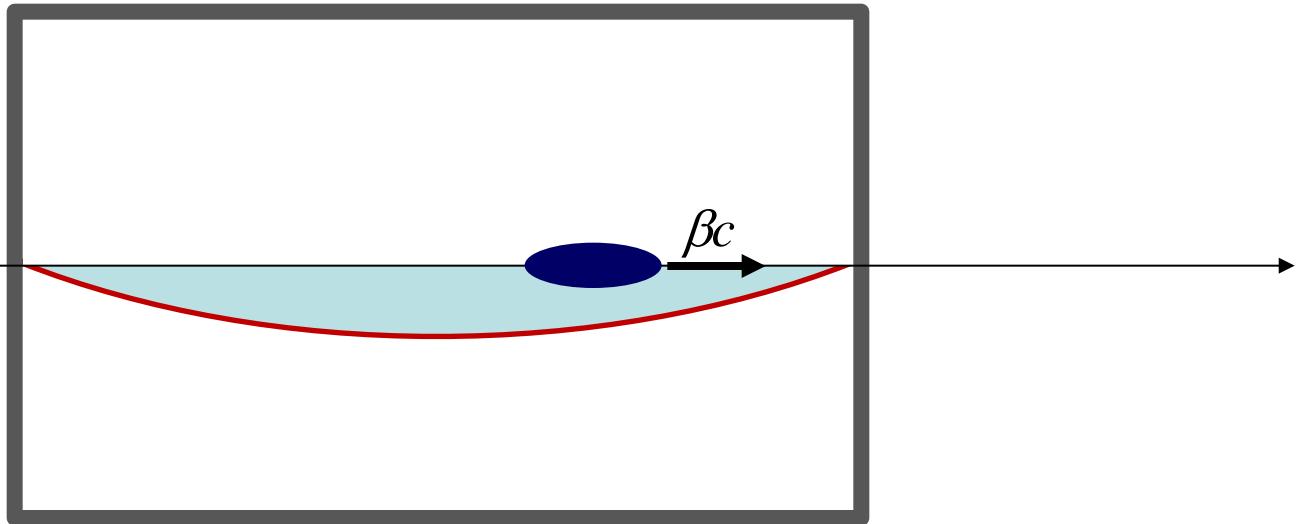
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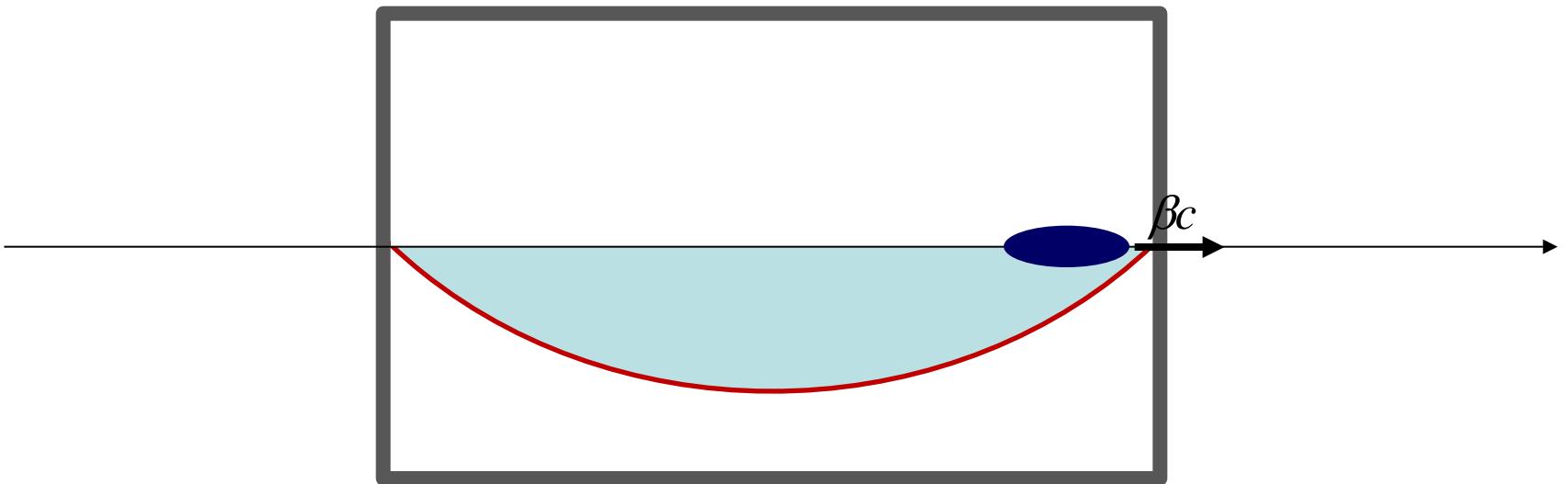
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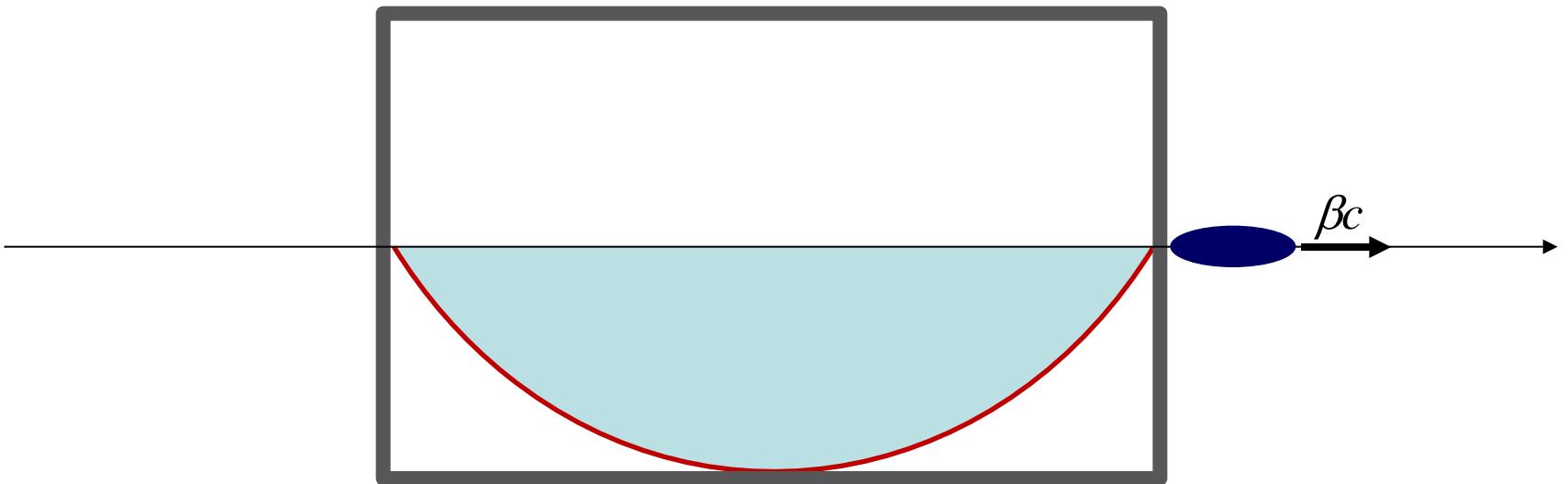
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Longitudinal Mode



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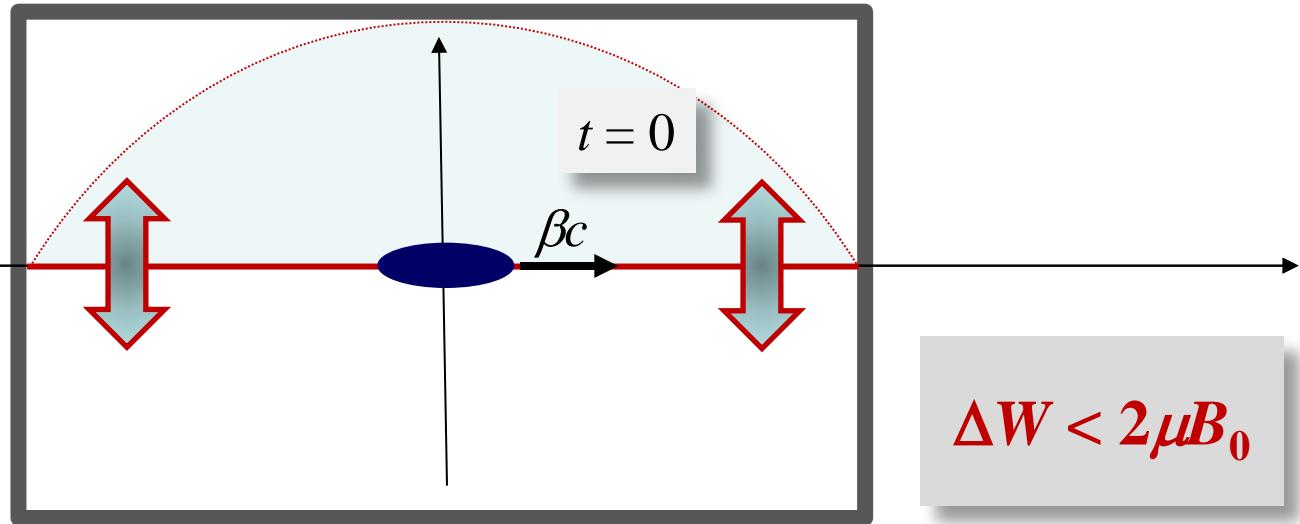
Longitudinal Mode



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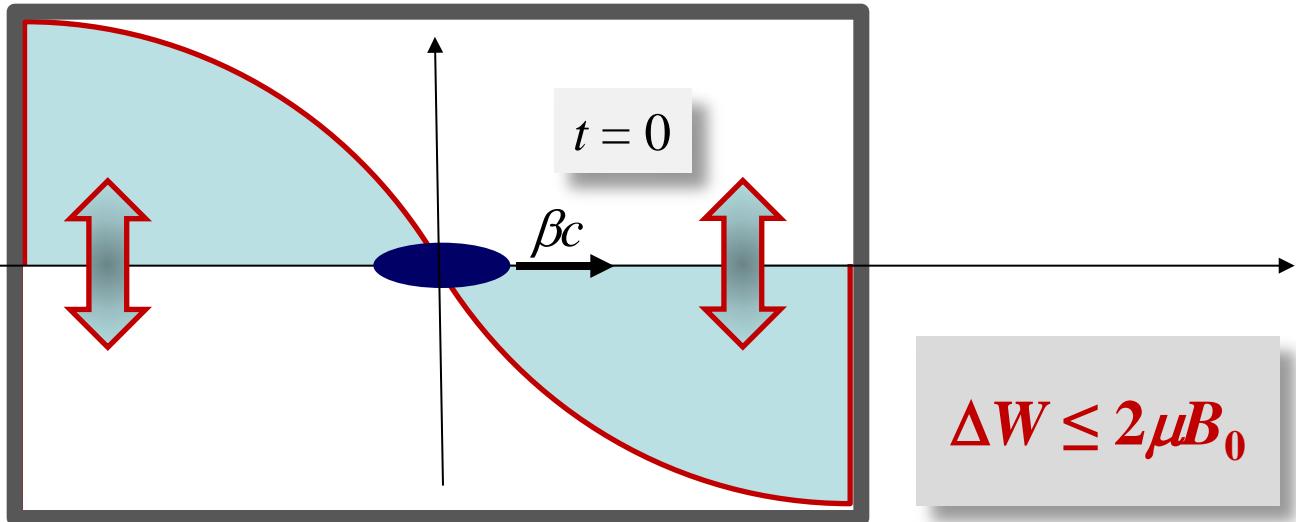
Findings:

$$B_{\perp} = B_0 \cdot \cos(\omega t + \phi) \Rightarrow \phi_{opt} = -\frac{\pi}{2}, \quad \beta_{ph} \approx 1$$

?

Transverse Mode

$$\Delta W = \int \frac{\partial}{\partial z} (\vec{\mu} \cdot \vec{B}) \cdot dz$$



Findings:

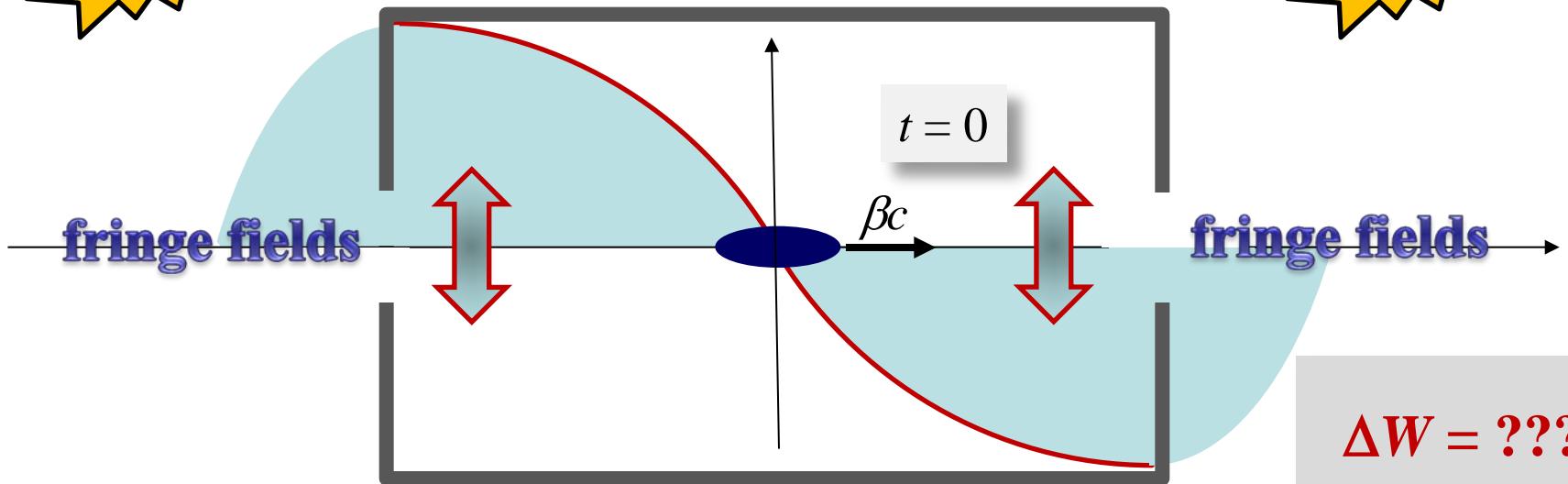
$$B_{\perp} = B_0 \cdot \cos(\omega t + \phi) \Rightarrow \phi_{opt} = 0, \beta_{ph} \gg 1$$

Transverse Mode

but:

$$\Delta W = \int \frac{\partial}{\partial z} (\vec{\mu} \cdot \vec{B}) \cdot dz$$

but:



$$\Delta W = ???$$

Findings:

$$B_{\perp} = B_0 \cdot \cos(\omega t + \phi) \Rightarrow \phi_{opt} = 0, \quad \beta_{ph} = ???$$

A simple but (hopefully) correct Approach

Transformation of derivatives: $\frac{\partial}{\partial z^*} = \gamma \left(\frac{\partial}{\partial z} + \frac{\beta}{c} \frac{\partial}{\partial t} \right) = \gamma \frac{d}{dz} - \frac{1}{\beta \gamma c} \frac{\partial}{\partial t}$

Transformation of the fields:

$$\vec{\mu}^* \cdot \vec{B}^* = \vec{\mu} \cdot \left[\frac{\gamma}{1+G} \left\{ \left(G + \frac{1}{\gamma} \right) \vec{B}_\perp - \left(G + \frac{1}{1+\gamma} \right) \frac{\vec{\beta}}{c} \times \vec{E} \right\} + \vec{B}_\parallel \right]$$

Taking use of the relativistic compensation:

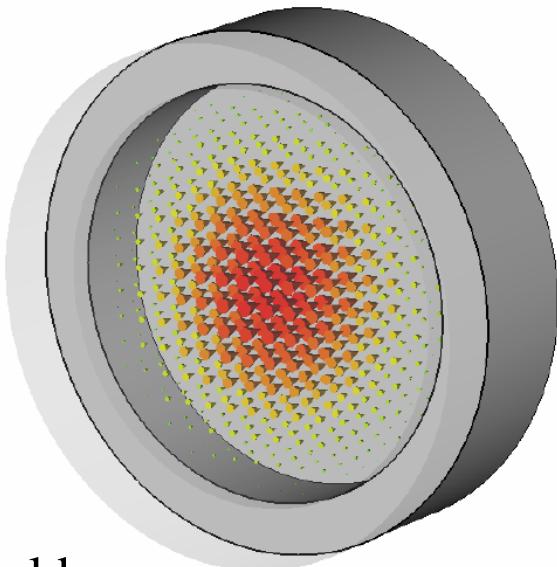
$$\Delta U = \int_0^d F_z^{SG} \cdot dz = \underbrace{\gamma \vec{\mu}^* \cdot \vec{B}^*}_{=0} \Big|_0^d - \frac{\vec{\mu}^*}{\beta c} \cdot \int_0^d \frac{\partial}{\partial t} \left[\frac{\gamma}{1+G} \left\{ \left(G + \frac{1}{\gamma} \right) \vec{B}_\perp - \left(G + \frac{1}{1+\gamma} \right) \frac{\vec{\beta}}{c} \times \vec{E} \right\} + \vec{B}_\parallel \right] dz$$

→ Energy transfer to the cavity:

$$\Delta U = \int_c^d F_z^{SG} \cdot dz = - \frac{\vec{\mu}}{\beta c} \cdot \frac{\partial}{\partial t} \int_c^d \left\{ \underbrace{\frac{G + \frac{1}{\gamma}}{1+G} \vec{B}_\perp}_{=\xi_B} - \underbrace{\left(\frac{G}{1+G} + \frac{1}{(1+G)(1+\gamma)} \right) \frac{\vec{\beta}}{c} \times \vec{E} + \frac{1}{\gamma} \vec{B}_\parallel}_{=\xi_E} \right\} dz$$

Cavity Modes: TM

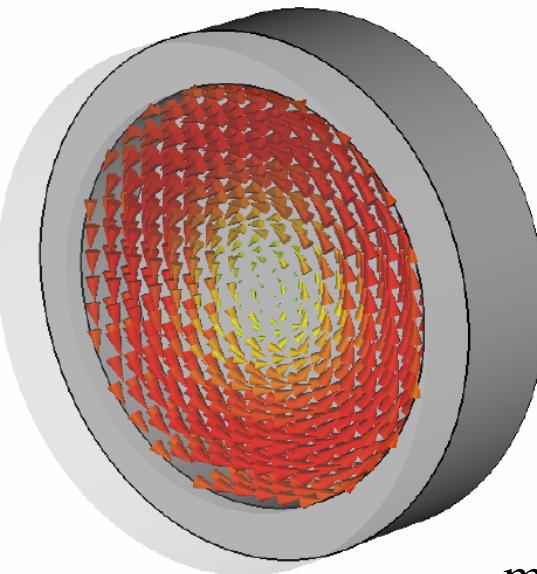
TM_{010}



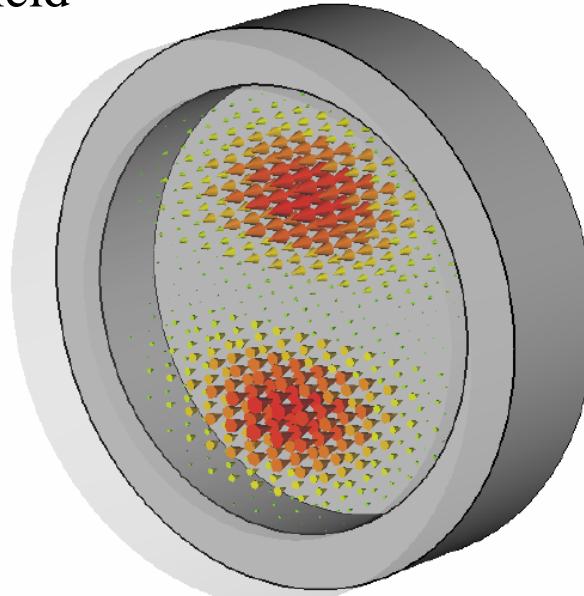
electric field



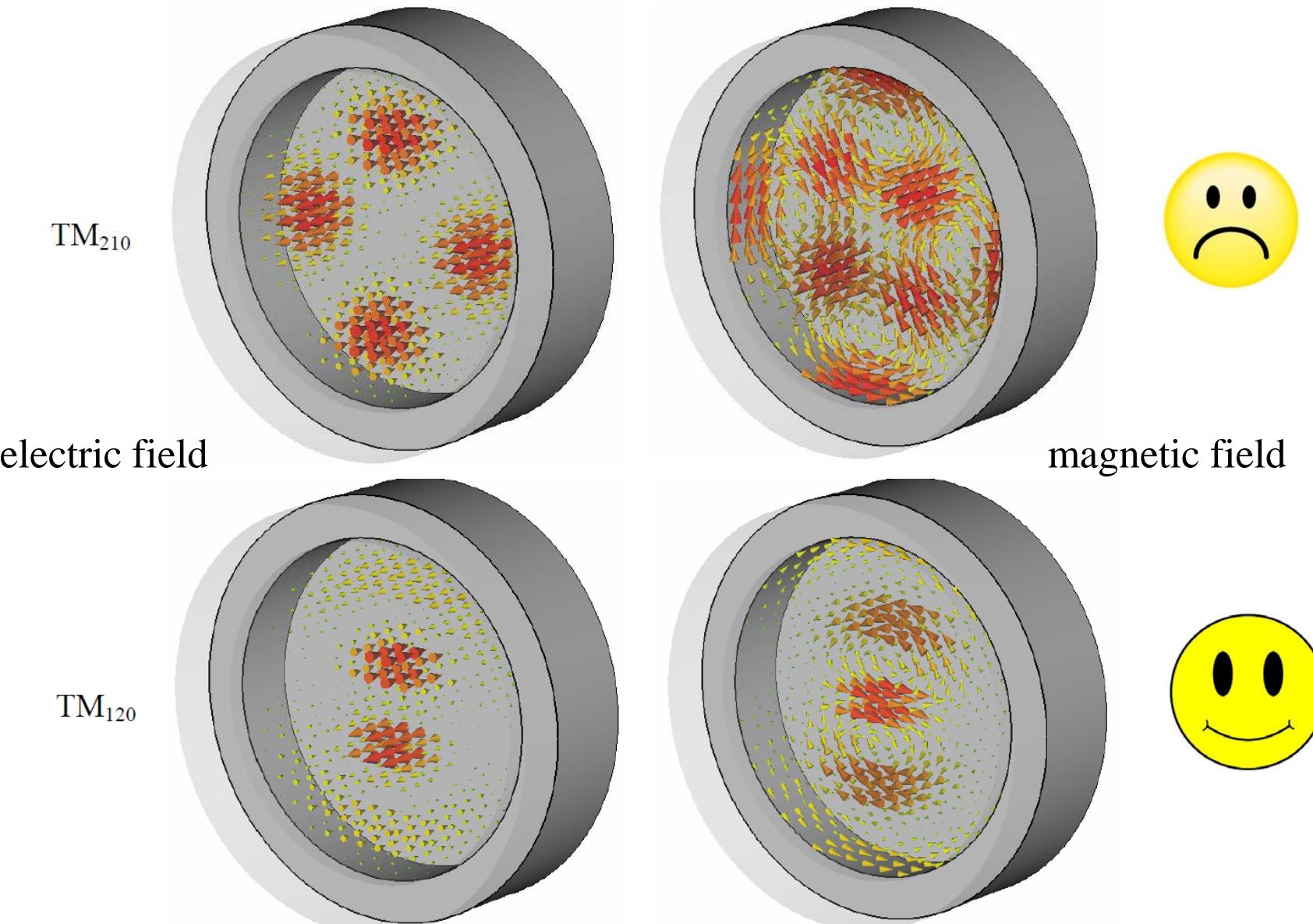
magnetic field



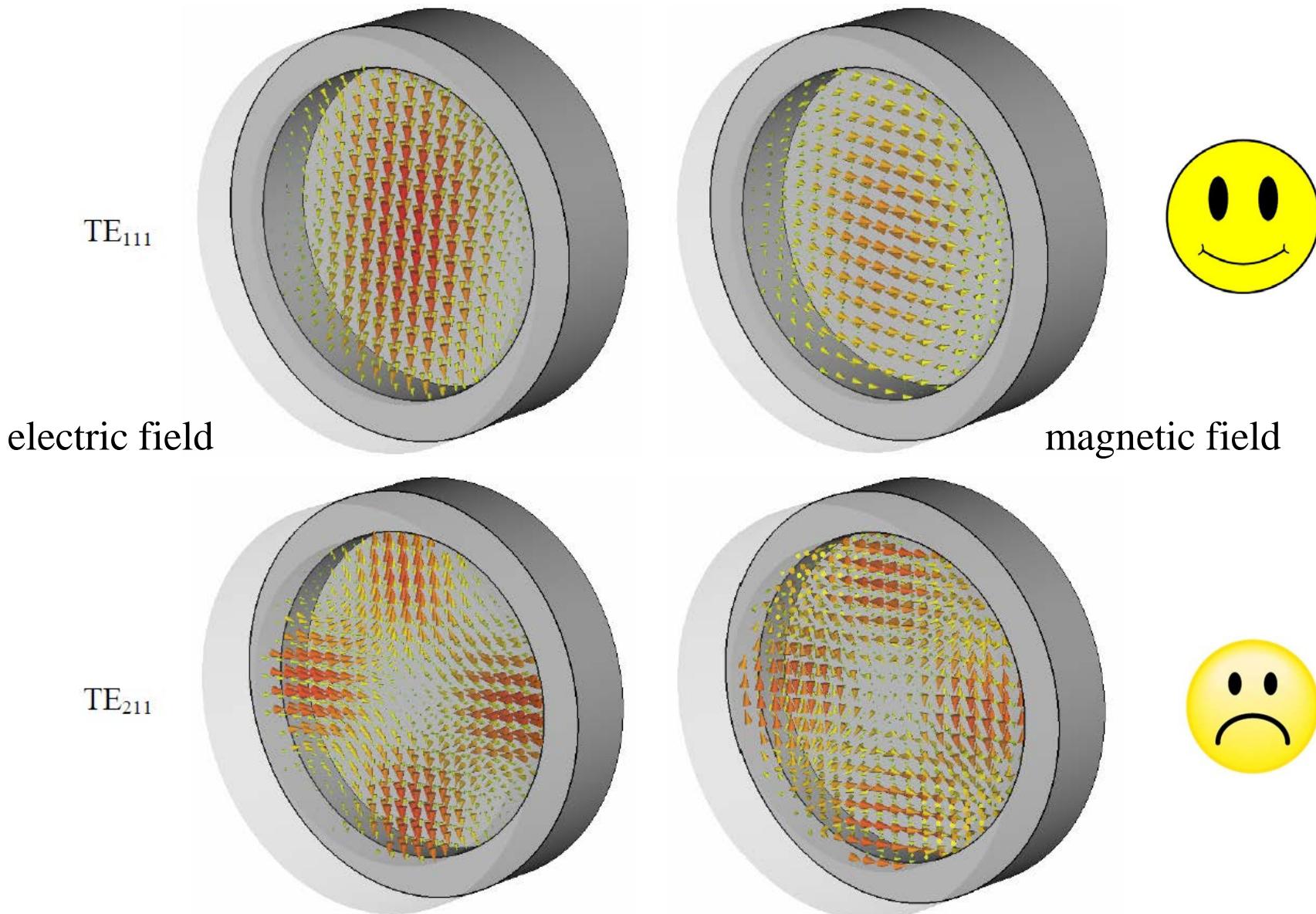
TM_{110}



Cavity Modes: TM



Cavity Modes: TE

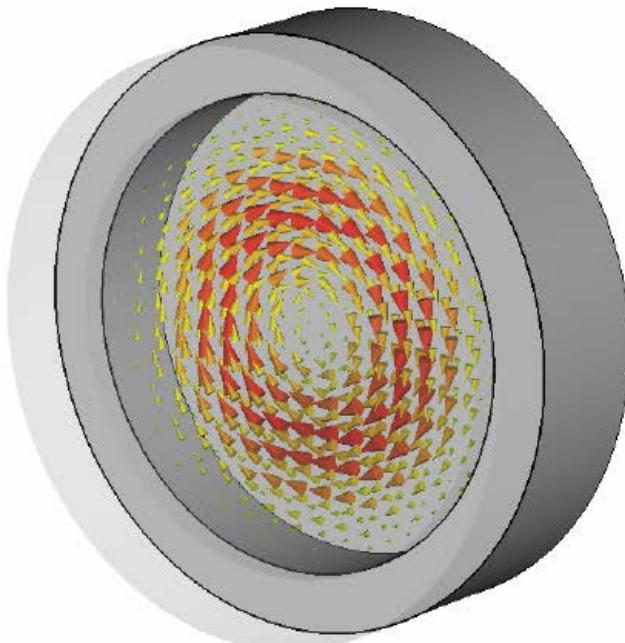


Cavity Modes: TE

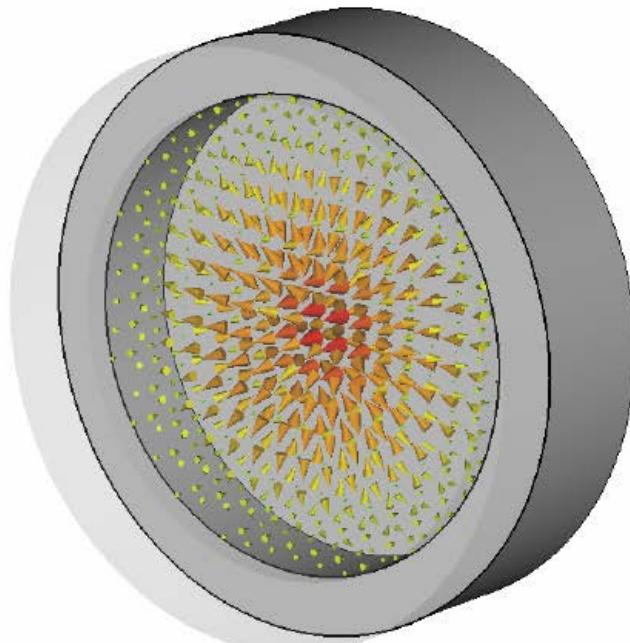
and longitudinal:

electric field

TE_{011}



magnetic field



Single Particle Energy Transfer

Integration of the Stern-Gerlach force:

- odd longitudinal p :

$$\Delta U_{\perp} = \frac{-2 \cos \phi}{1 - (\beta/\beta_{ph})^2} \sin\left(\frac{p\pi}{2}\right) \cos\left(\frac{p\pi\beta_{ph}}{2\beta}\right) \vec{\mu} \cdot \left\{ \xi_B \vec{B}_{\perp}^0 + \xi_E \frac{\beta}{\beta_{ph}} \left(\hat{e}_z \times \frac{\beta}{c} \vec{E}_{\perp}^0 \right) \right\}$$

$$\Delta U_{\parallel} = -\frac{2}{\gamma} \mu_z B_z^0 \frac{\sin \phi}{1 - (\beta/\beta_{ph})^2} \frac{\beta}{\beta_{ph}} \sin\left(\frac{p\pi}{2}\right) \cos\left(\frac{p\pi\beta_{ph}}{2\beta}\right)$$

- even longitudinal p :

$$\Delta U_{\perp} = \frac{2 \sin \phi}{1 - (\beta/\beta_{ph})^2} \cos\left(\frac{p\pi}{2}\right) \sin\left(\frac{p\pi\beta_{ph}}{2\beta}\right) \vec{\mu} \cdot \left\{ \xi_B \vec{B}_{\perp}^0 - \xi_E \frac{\beta}{\beta_{ph}} \left(\hat{e}_z \times \frac{\beta}{c} \vec{E}_{\perp}^0 \right) \right\}$$

$$\Delta U_{\parallel} = \frac{2}{\gamma} \mu_z B_z^0 \frac{\cos \phi}{1 - (\beta/\beta_{ph})^2} \frac{\beta}{\beta_{ph}} \cos\left(\frac{p\pi}{2}\right) \sin\left(\frac{p\pi\beta_{ph}}{2\beta}\right)$$

Signal Power

Energy transfer: $P_+ = \frac{I}{e} \cdot \Delta U$, **bunch factor:** $\eta_b = \int \rho(s) \cdot \cos\left(\frac{\omega s}{\beta c}\right) \cdot ds$

Stored energy: $W_C = \frac{1}{2\mu_0} \int_V B^2 dV = \frac{1}{2\varepsilon_0} \int_V E^2 dV = v_b \cdot B_0^2 = v_e \cdot E_0^2$

→ **Energy transfer:** $dW_C = P_+ \cdot dt = \frac{I}{e} \cdot \eta_b \cdot \Delta U \cdot dt = \frac{I}{e} \cdot \eta_b \cdot s_\mu \cdot B_0 \cdot dt = \varsigma \cdot \sqrt{W_C} \cdot dt$

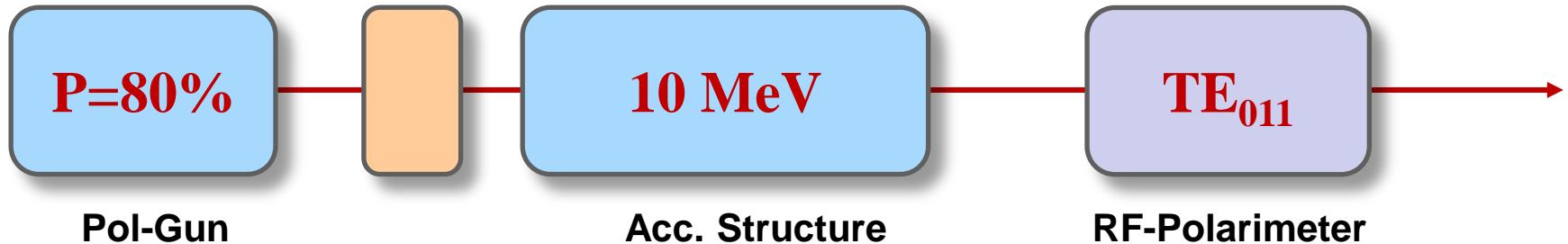
Energy dissipation: $P_- = \frac{\omega}{Q_l} \cdot W_C = \frac{1+\kappa}{Q_0} \cdot \omega \cdot W_C = \frac{1}{\tau} \cdot W_C$

Build-up of stored energy: $\frac{d}{dt} W_C = \varsigma \cdot \sqrt{W_C} - \frac{1}{\tau} \cdot W_C \quad \rightarrow \quad W_C(t) = (\varsigma\tau)^2 \cdot \left(1 - e^{-\frac{t}{2\tau}}\right)$

Steady state conditions: $W_C^\infty = (\varsigma\tau)^2 = \frac{I^2 \cdot \eta_b^2 \cdot s_\mu^2}{e^2 \cdot v} \cdot \frac{Q_0^2}{(1+\kappa)^2} \cdot \frac{1}{\omega^2}$

Signal Power: $P_s = \kappa \cdot P_-^C = \kappa \cdot \frac{\omega \cdot W_C}{Q_0} = \frac{I^2 \cdot \eta_b^2 \cdot s_\mu^2}{e^2 \cdot v} \cdot \frac{\kappa}{(1+\kappa)^2} \cdot \frac{Q_0}{\omega}$

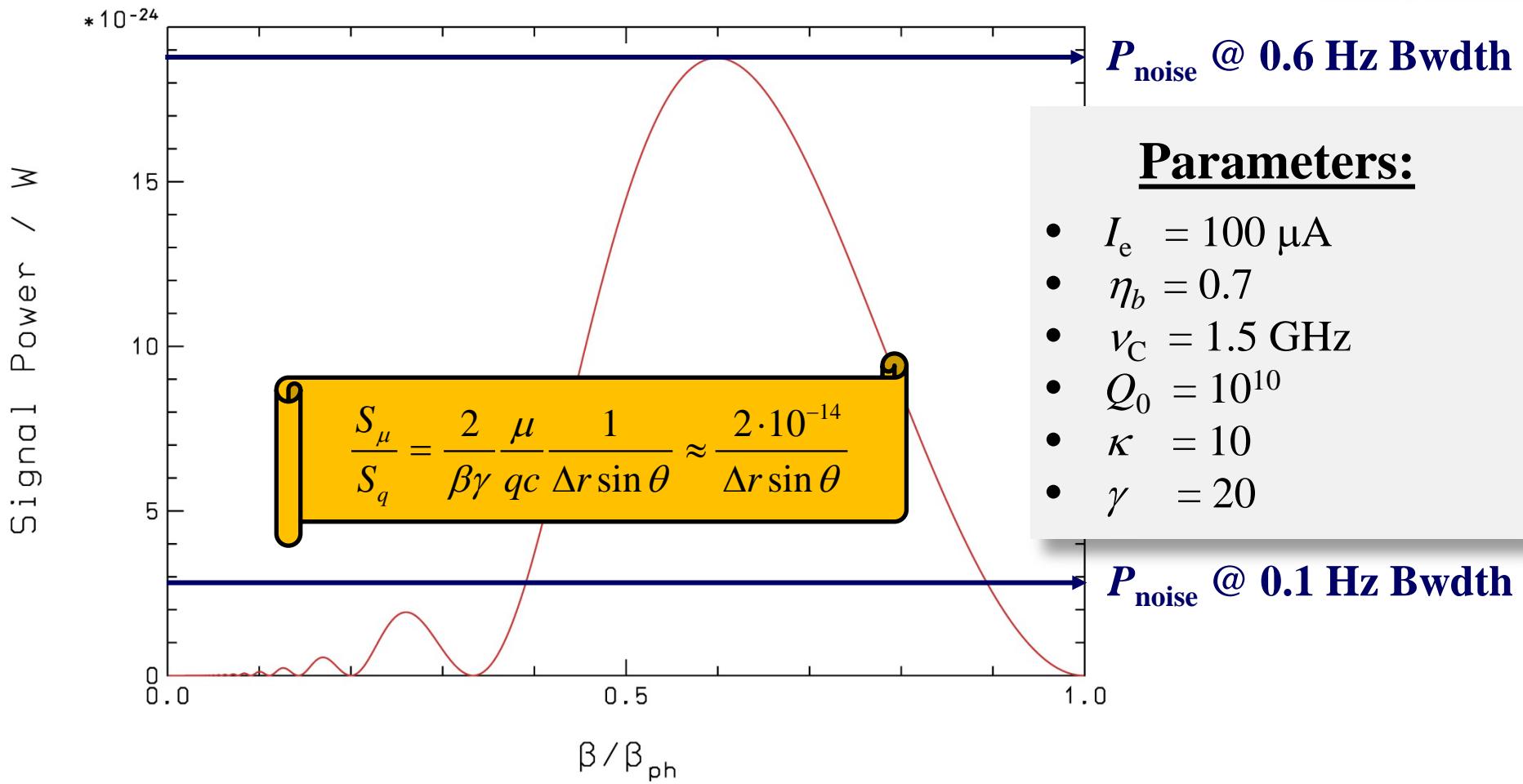
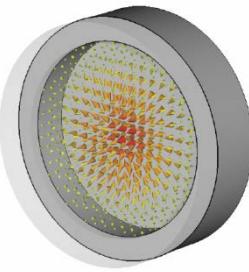
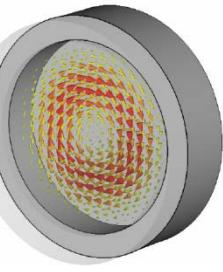
Experiment @ JLAB:



PoP Test at the injector:

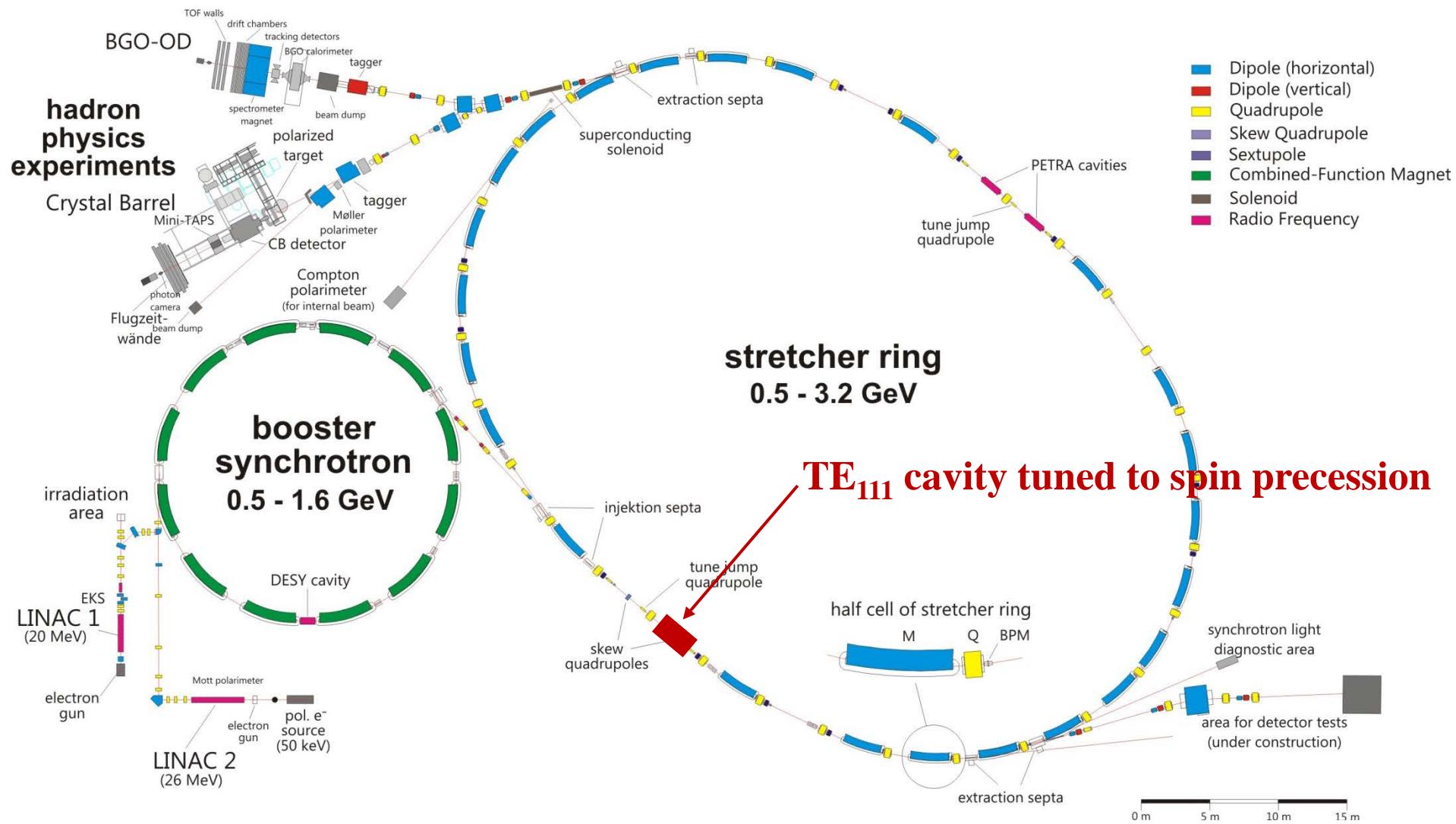
- Longitudinal polarisation \leftrightarrow long. magn. field
- Low Lorentz gamma
- Flip helicity with Pockels cell
- Tune cavity to bunch repetition frequency
- Use TE mode with no long. electric fields
- Phase locking of polarimeter signal to RF

Longitudinal: TE₀₁₁

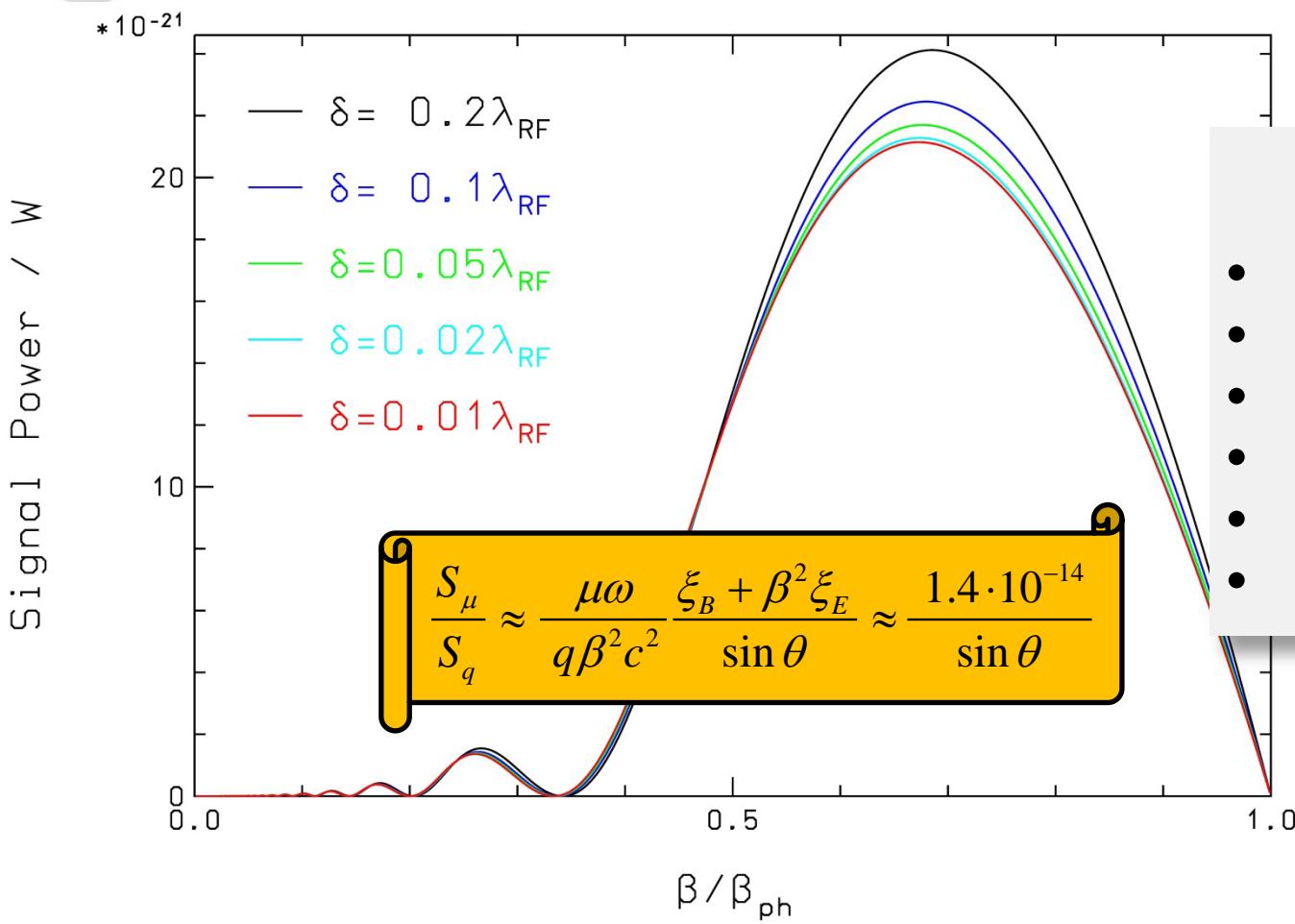
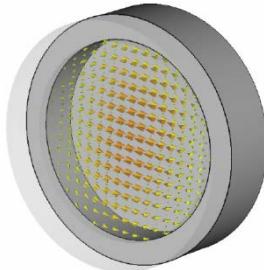
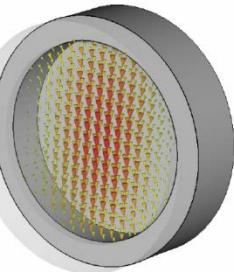


Expected Signal Power: $P_s = \left(\frac{I \cdot \eta_b}{e} \right)^2 \cdot \frac{16\mu_0\mu_e^2}{\pi^2 c^3} \cdot \frac{f(\beta_{ph})}{F(j_{11})} \cdot \frac{\kappa Q_0}{(1+\kappa)^2} \cdot \left(\frac{\omega_c}{\gamma} \right)^2$

Experiment @ ELSA



Transverse: TE₁₁₁



Parameters:

- $I_e = 50 \text{ mA}$
- $\eta_b = 0.7$
- $\nu_C = 1.5 \text{ GHz}$
- $Q_0 = 10^{10}$
- $\kappa = 10$
- $\gamma = 2000$

Expected Signal Power: $P_s \approx \left(\frac{I \cdot \eta_b}{e}\right)^2 \cdot \frac{32\mu_0\mu_e^2}{\pi^2 c^3} \cdot \frac{f(\beta_{ph})}{F(j'_{11})} \cdot \frac{\kappa Q_0}{(1+\kappa)^2} \cdot (\mathbf{G} \cdot \boldsymbol{\omega}_c)^2$

Conclusions

- Expected signal power is extremely low!
- sc cavities ($Q_0 \approx 10^{10}$) with weak coupling essential!
- Phase-lock techniques required
- Coupling to charge is about 14 orders of magnitude greater!

PoP will be a really hard task but doable??

LIGO demonstrated: ultimate precision can be achieved!

Stern-Gerlach

May the force be with us!

