

Exercise Sheet 3, Theoretical Quantum and Atom Optics

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Exercise 5. Coherent States

The coherent states $\{|\alpha\rangle, \alpha \in \mathbb{C}\}$ are an overcomplete non-orthogonal “basis“ of the Fock space:

- (a) Show that the orthogonality relation $\langle\alpha|\alpha'\rangle = \delta(\alpha - \alpha')$ does not hold for coherent states by explicitly computing the scalar product $\langle\alpha|\alpha'\rangle$.
- (b) Prove the following completeness relation of the coherent states:

$$\int d^2\alpha |\alpha\rangle \langle\alpha| = \pi.$$

- (c) Combining the above results, show that for a given coherent state $|\beta\rangle$ there are at least two different ways of expanding it in terms of the coherent states $\{|\alpha\rangle, \alpha \in \mathbb{C}\}$.

Hints:

For the calculations expand the coherent state in Fock states $|\alpha\rangle = \sum_n |n\rangle \langle n|\alpha\rangle$.

Use $\alpha = ae^{i\phi}$ and $d^2\alpha = adad\phi$ before performing the integration and use the definition of the gamma function to solve the integral.

5 Points

Exercise 6. Squeeze operator, squeezed states

Consider the squeeze operator

$$S(\xi) = \exp\left(\frac{1}{2}\xi^* a^2 - \frac{1}{2}\xi(a^\dagger)^2\right),$$

where $\xi = re^{i\theta}$ is a complex number.

- (a) Show that $S^\dagger(\xi) = S(-\xi)$. Further conclude that $S(-\xi) = S^{-1}(\xi)$ by showing that $S(\xi)S(-\xi) = 1$.
- (b) Now consider the squeezed coherent state $|\alpha, \xi\rangle = S(\xi)D(\alpha)|0\rangle$. Explicitly calculate the expectation values $\langle a \rangle = \langle\alpha, \xi|a|\alpha, \xi\rangle$ and $\langle a^2 \rangle = \langle(\hat{a}^\dagger)^2\rangle^*$.
- (c) Y_1 and Y_2 are the rotated quadrature variables, see the lecture. Explicitly calculate the variances $(\Delta Y_1)^2$ and $(\Delta Y_2)^2$ and show that $\Delta Y_1 \Delta Y_2 = \frac{1}{4}$.
- (d) There are (at least) two ways for creating a coherent state $|\alpha\rangle$ out of the vacuum $|0\rangle$. One can either employ the unitary displacement operator, i.e. $|\alpha\rangle = \hat{D}(\alpha)|0\rangle$, or one can write $|\alpha\rangle = \tilde{N} e^{\alpha\hat{a}^\dagger}|0\rangle$ with $\tilde{N} = 1/||e^{\alpha\hat{a}^\dagger}|0\rangle||$ (which is fun to prove). Similarly, one can generate the squeezed state $|\xi\rangle = \hat{S}(\xi)|0\rangle$ by means of the simpler operator $e^{-\zeta/2(\hat{a}^\dagger)^2}$ as follows: The choice $\zeta = \tanh(|\xi|)\frac{\xi}{|\xi|}$ leads to $|\xi\rangle = N e^{-\zeta/2(\hat{a}^\dagger)^2}|0\rangle$ with $N = 1/||e^{-\zeta/2(\hat{a}^\dagger)^2}|0\rangle||$.

Consider the Hamiltonian $\hat{H} = \omega(\hat{a}^\dagger\hat{a} + \frac{1}{2})$ and a squeezed state as the initial condition, i.e. $|\psi(t=0)\rangle = |\xi\rangle$. You are now asked to solve the time-dependent Schrödinger equation and to interpret your solution $|\psi(t)\rangle$.

5 Points