## 12. Excercise Sheet, Lecture Theoretical Quantum and Atom Optics

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To be returned on 09/07/2012, in the tutorials

## Exercise 24. Vortices in Bose-Einstein Condensates

Let  $\psi(\mathbf{r})$  denote the Gross-Pitaevskii wavefunction of a repulsively interacting condensate in free space. Certainly,  $\psi(\mathbf{r})$  should be a single-valued quantity as any wavefunction. In particular, the value of  $\psi(\mathbf{r})$  must coincide with the value of the wavefunction if one starts at  $\mathbf{r}$  and returns to  $\mathbf{r}$  after moving along any closed loop  $\zeta$ .

(a) Consider the Gross-Pitaevskii wavefunction in its hydrodynamic representation, i.e.  $\psi(\mathbf{r}) = f(\mathbf{r}) e^{i\phi(\mathbf{r})}$ . Show that for any closed loop  $\zeta$  the circulation:

$$C_{\zeta}[\boldsymbol{v}] = \oint_{\zeta} d\boldsymbol{r} \, \boldsymbol{v}(\boldsymbol{r}) \tag{1}$$

is quantized, where  $\mathbf{v}(\mathbf{r}) = \frac{\hbar}{m} \nabla \phi(\mathbf{r})$  denotes the local velocity of the condensate. Quantization of the circulation means that there is a constant  $\alpha$  such that  $C_{\zeta}[\mathbf{v}] = \alpha l$  for some  $l \in \mathbb{Z}$ . Determine  $\alpha$ .

- (b) Now show that applying Stokes' theorem to equation (1) in a naive way would lead to the conclusion that l=0 always. This, however, contradicts the experimental observations that BECs can feature vortices<sup>1</sup> and even turbulence! So why is Stokes' theorem not applicable in this situation? Please be reminded of the exercise sheet 4 ("the phase operator"). There, one central result was that it is only possible to define an approximately hermitian phase operator if the probability of having a zero density vanishes. So how does the density profile of a vortex have to look like qualitatively?
- (c) Now let's turn to a 2-dimensional BEC living in the xy-plane. Suppose that there is a vortex of so-called *charge* l at the origin. Introducing polar coordinates  $(\rho, \varphi)$ , assume that the phase field  $\phi(\mathbf{r})$  is independent of  $\rho$ . Make the simplest ansatz for  $\phi(\mathbf{r})$  leading to a vortex of charge l and determine the corresponding velocity field  $\mathbf{v}(\mathbf{r})$ !
- (d) Finally, we aim at calculating the energy of a single vortex of charge l, which is defined as the difference of the energy of a uniform BEC with and without vortex. For obtaining a finite energy, we should better consider a disk of radius R out of the uniform condensate, where the vortex shall lie in the origin of this disk. Show that the energy functional for the uniform condensate with the vortex equals:

$$E_l[f] = 2\pi \int_{0}^{R} d\rho \, \rho \left[ \frac{\hbar^2}{2m} (\partial_\rho f)^2 + \frac{\hbar^2}{2m} l^2 \frac{f^2}{\rho^2} + \frac{U_0}{2} f^4 \right], \tag{2}$$

where an isotropic density distribution  $f(\rho)$  has been assumed. The density profile can be determined by minimizing  $E_l[f]$  under the constraint that  $\int d\rho \, \rho f(\rho)$  equals the total particle number. The resulting ODE, however, can only be handled numerically. In contrast to this, the energy of the BEC without vortex,  $E_{l=0}[f_0]$ , can be easily calculated since  $f_0$  is a constant. One can show that for  $R >> \xi = \sqrt{\hbar^2/2mU_0f_0^2}$  ( $\xi$  is the so called healing length) the energy of the singly charged vortex is given by:

$$\varepsilon_l \equiv E_l[f] - E_{l=0}[f_0] \approx \pi \frac{\hbar^2}{m} f_0^2 l^2 ln \frac{R}{|l|\xi}$$
(3)

Would it be energetically favourable to have a single vortex of (not too large) charge l > 0 or l spatially well separated singly charged vortices?

10 Points

<sup>&</sup>lt;sup>1</sup>A BEC is said to have at least one vortex within the boundary  $\zeta$  if  $C_{\zeta}[v] \neq 0$ .