

Winter term 2012/13  
**Exercise Sheet 11, Theoretical Quantum and Atom Optics**  
 University of Hamburg, Prof. P. Schmelcher

To be returned on Tuesday, 22/01/2013, in the tutorials

**Exercise 21.** Collective excitations of Bose-Einstein condensates

Consider a repulsively interacting Bose-Einstein condensate ( $U_0 > 0$ ) in an isotropic 3-dimensional harmonic trap:  $V(\mathbf{r}) = \frac{m\omega_0^2}{2}r^2$ . Furthermore, let us assume that the number of atoms is so large that we can apply the Thomas-Fermi approximation leading to the ground state density profile:

$$n_0(\mathbf{r}) = \frac{\mu - V(\mathbf{r})}{U_0} \quad \text{for } r \leq R = \left(\frac{2\mu}{m\omega_0^2}\right)^{\frac{1}{2}}. \quad (1)$$

In the following, we aim at discussing collective excitations of the atom cloud, for which the hydrodynamic formulation of the Gross-Pitaevskii theory turns out to be a convenient approach (cf. sheet 8). With  $\psi = \sqrt{n(\mathbf{r}, t)}e^{i\varphi(\mathbf{r}, t)}$ , we have obtained:

$$\frac{\partial n}{\partial t} = -\nabla \cdot (n\mathbf{v}), \quad (2)$$

$$-\hbar \frac{\partial \varphi}{\partial t} = -\frac{\hbar^2}{2m\sqrt{n}} \Delta \sqrt{n} + \frac{1}{2}m\mathbf{v}^2 + V + nU_0, \quad (3)$$

where the velocity field is given as  $\mathbf{v} = \frac{\hbar}{m} \nabla \varphi$ .

- (a) Why can we neglect the quantum pressure term  $\propto (\Delta \sqrt{n})/\sqrt{n}$  here?
- (b) Show that  $\psi_0 = \sqrt{n_0(\mathbf{r})} \exp(-i\mu t/\hbar)$  solves the Eqs. (2,3) within this approximation.
- (c) Consider a small perturbation of this stationary solution,  $n(\mathbf{r}, t) = n_0(\mathbf{r}) + \delta n(\mathbf{r}, t)$ ,  $\varphi(\mathbf{r}, t) = -\mu t/\hbar + \delta \varphi(\mathbf{r}, t)$ , and linearize Eqs. (2,3) around the stationary solution. Show that the density fluctuation obeys:

$$\partial_t^2 \delta n = \frac{U_0}{m} ((\nabla n_0) \cdot (\nabla \delta n) + n_0 \Delta \delta n). \quad (4)$$

- (d) Make a separation ansatz  $\delta n = f(\mathbf{r}) e^{-i\omega t}$  and show that  $f$  satisfies:

$$\omega^2 f = \frac{1}{m} ((\nabla V) \cdot (\nabla f) - (\mu - V) \Delta f). \quad (5)$$

Explicate this equation for the given harmonic trap.

- (e) Introduce spherical coordinates  $x = r \sin \theta \cos \phi$ ,  $y = r \sin \theta \sin \phi$ ,  $z = r \cos \theta$ . Check that Eq. (5) is solved by the ansatz

$$f = Cr^l Y_{lm}(\theta, \phi), \quad C \in \mathbb{R}, l \in \mathbb{N}, m \in \{-l, -l+1, \dots, l\}. \quad (6)$$

where  $Y_{lm}$  denote the spherical harmonics, and find the corresponding eigenfrequency  $\omega$ . Look at this solution in more detail for parameters  $l = 1, m = 0$ . Show that for an infinitesimally small amplitude of the perturbation, the real part of the time-dependent total density can be written as  $\Re[n(\mathbf{r}, t)] = n_0(\mathbf{r} - \epsilon \cos(\omega_0 t) \mathbf{e}_z)$  with a (small) constant  $\epsilon$ . What does this oscillation look like?

- (f) Check that the ansatz  $f = C \left(1 - \frac{5}{3} \frac{r^2}{R^2}\right)$  also solves Eq. (5). What does the corresponding oscillation of the condensate density look like?

**10 Points**